



FIR/submm Astrophysics – An Overview

2 May 2018

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SOFIA Workshop

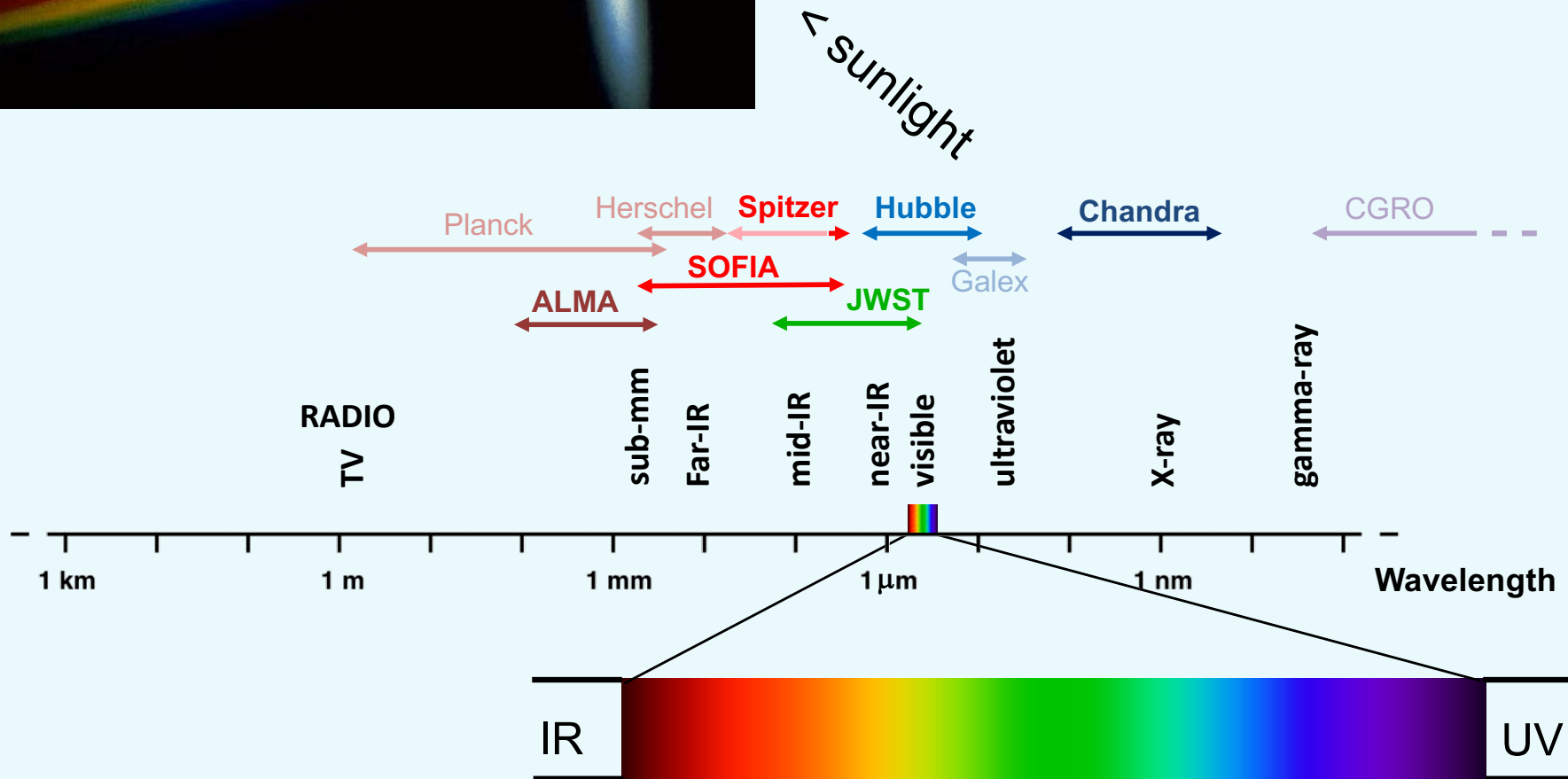


FIR – submm Astrophysics 2-4 May 2018



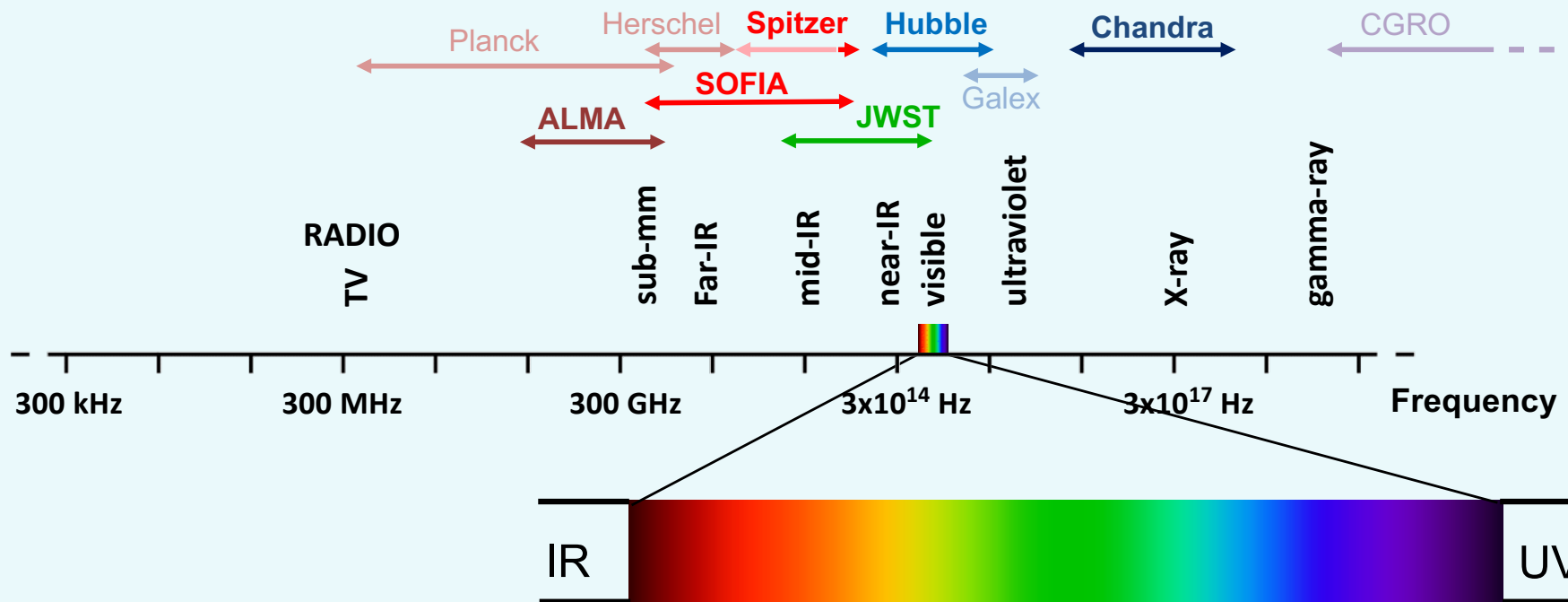


The Electromagnetic Spectrum



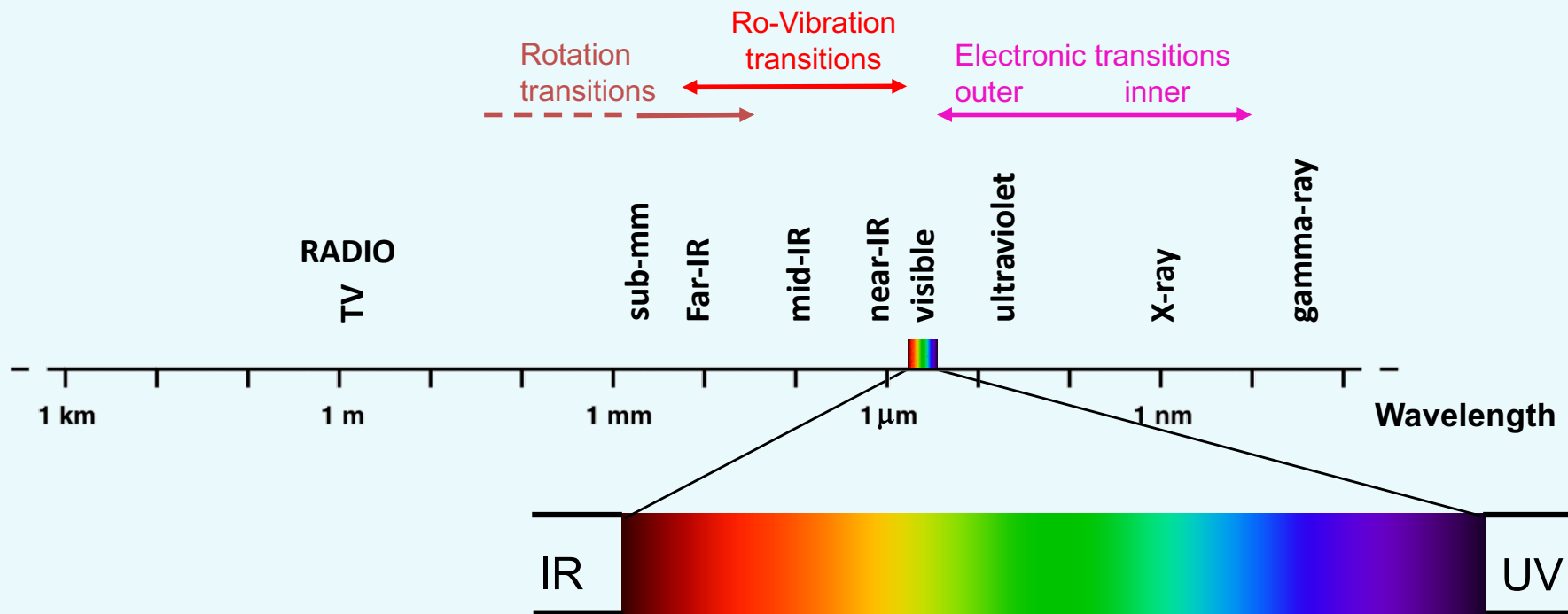


The Electromagnetic Spectrum

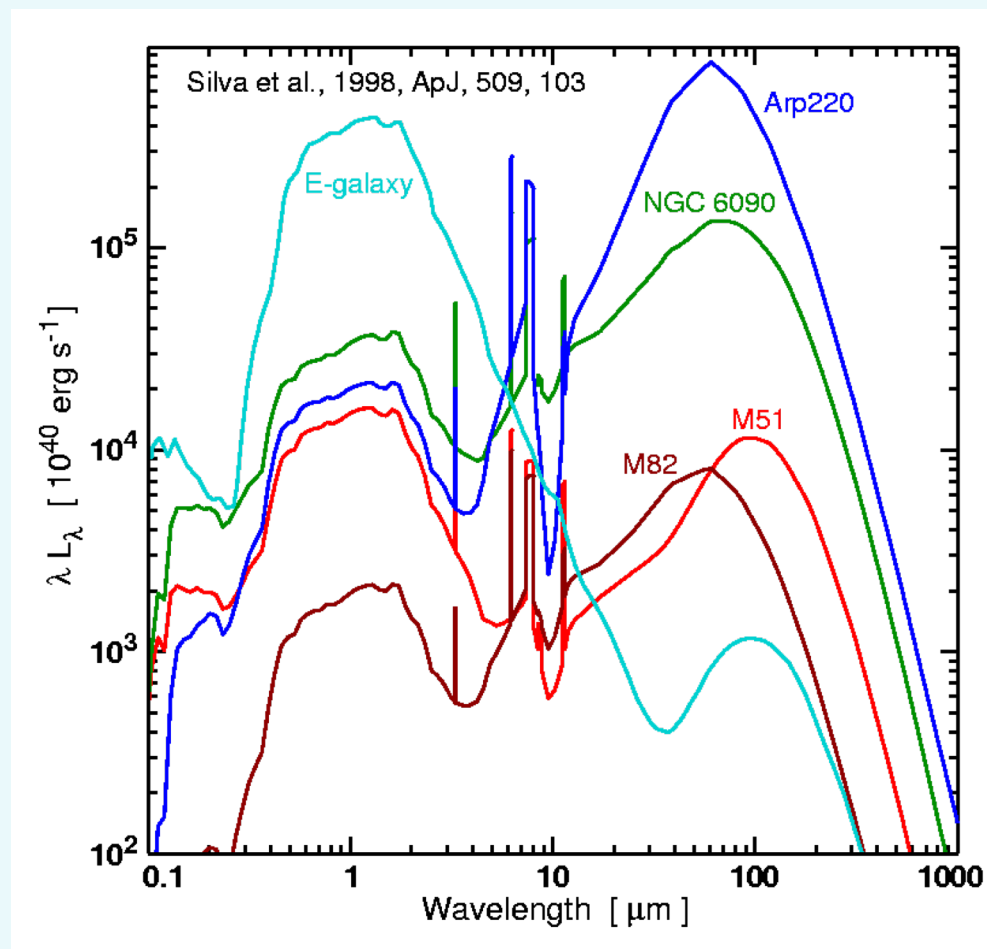




← sunlight



- Dusty galaxies emit mostly in the Far-IR. These wavelengths probe their star formation properties and evolution
- 1/2 of energy emitted since the Big Bang is in the Far-IR/submm
- Far-IR flux measures star formation activity and/or AGN activity
- [CII] 158 μ m is primary coolant of ISM gas, comprising $\sim 1\%$ of the Milky Way's total energy output. Other coolants and other fine structure lines provide important diagnostic information





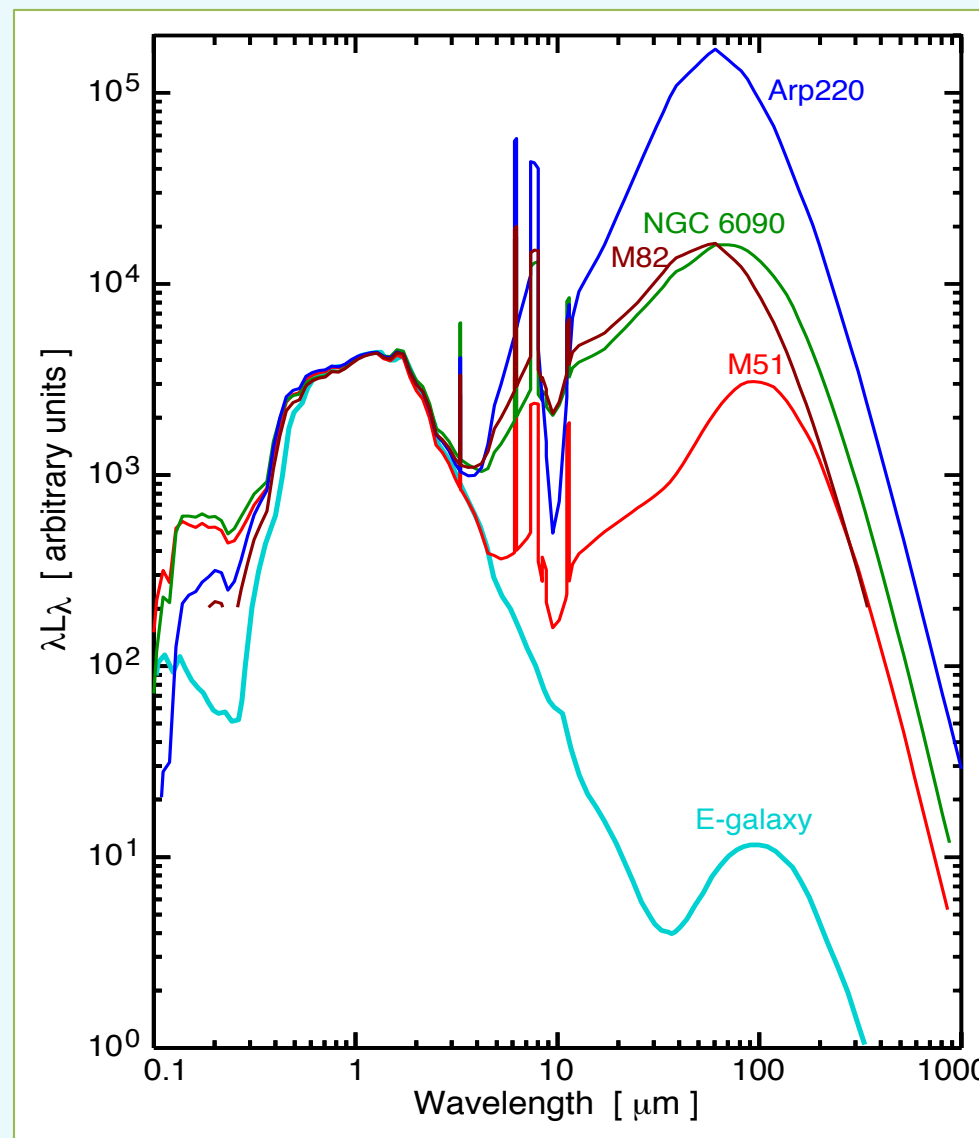
What Role does Dust play in the Structure and Evolution of the Universe?

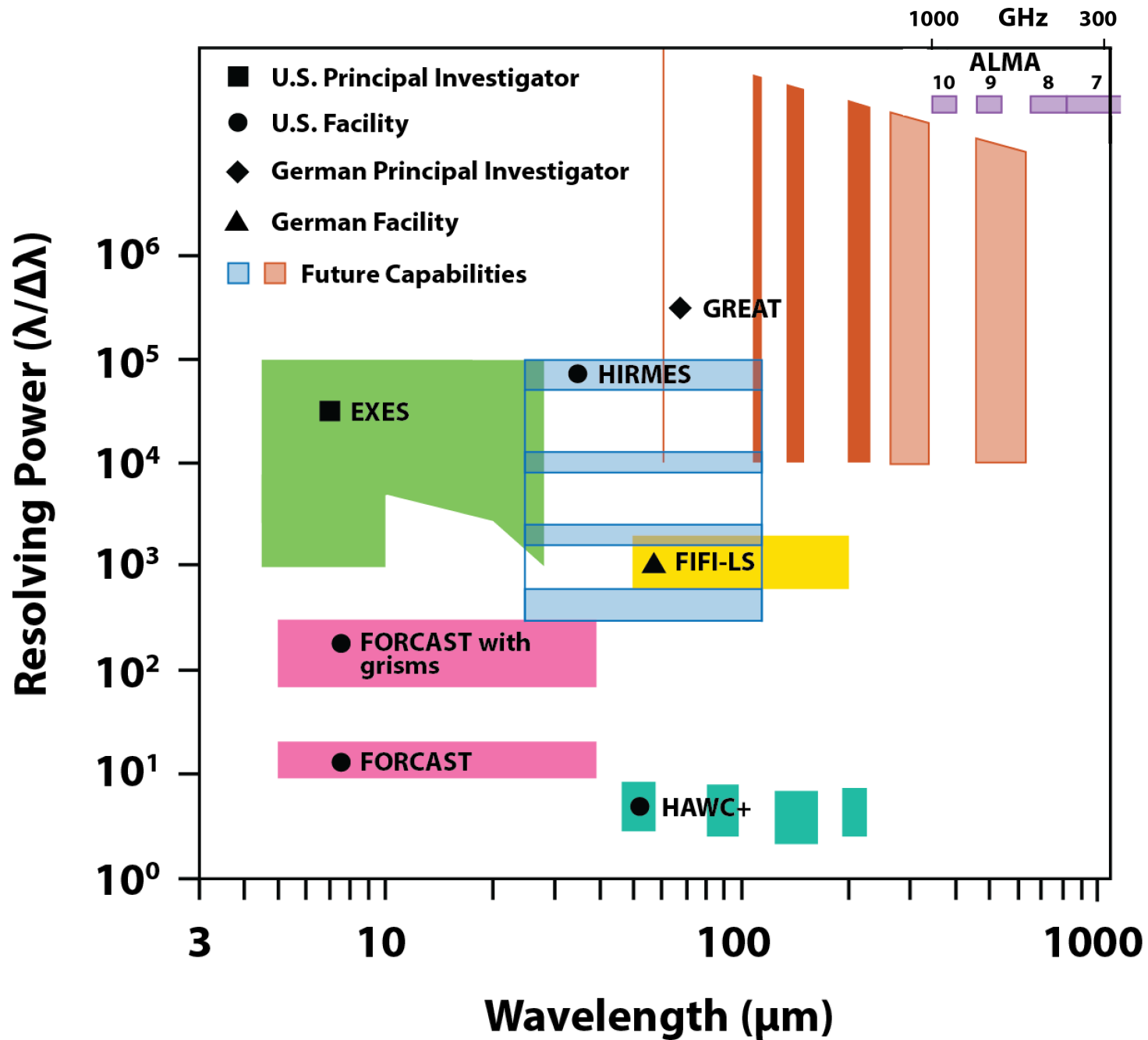


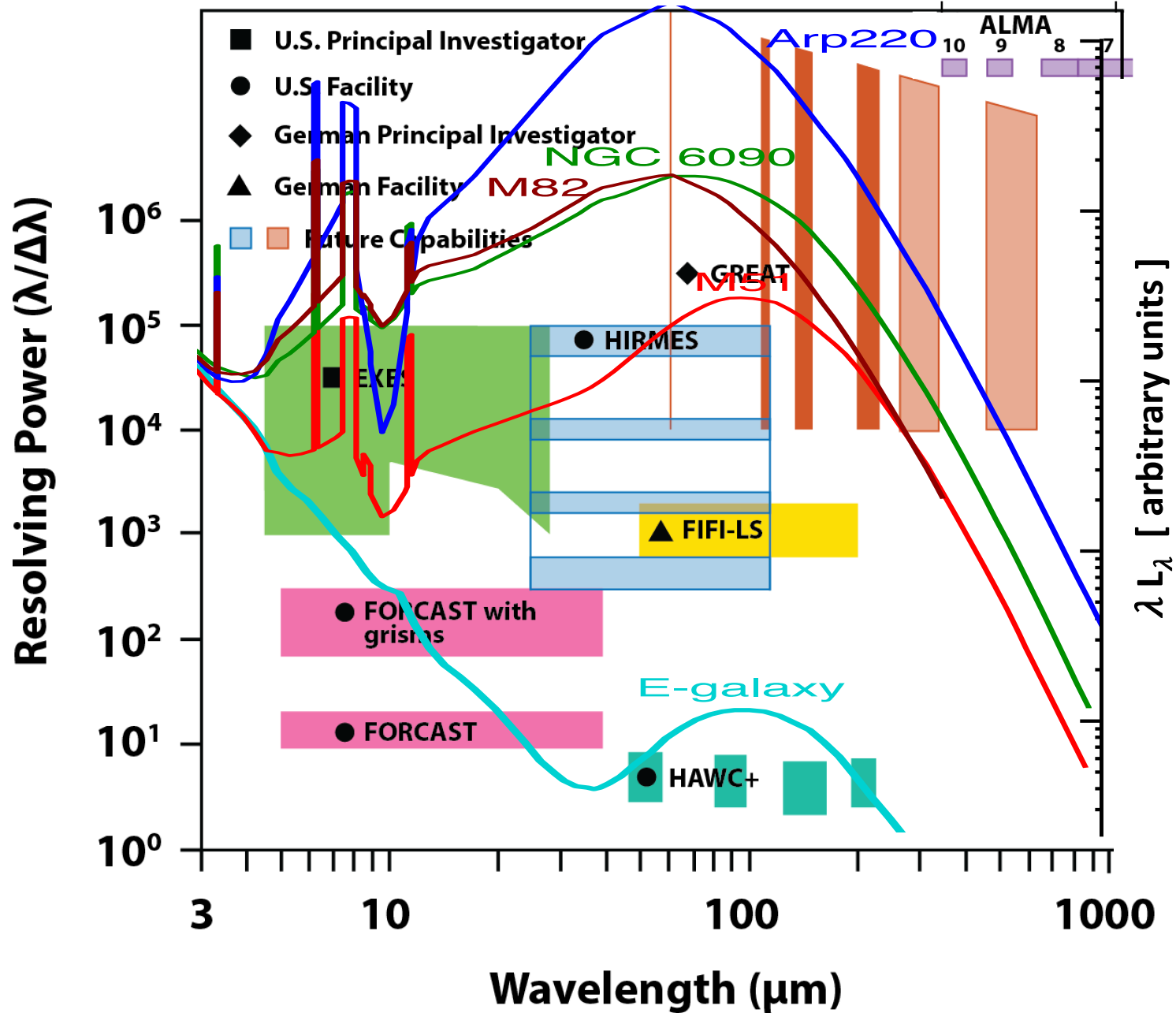
After the Big Bang, dust appears simultaneously with the first galaxies. Since then, it has played a crucial role in heating/cooling, molecule formation, and dynamics of galactic gas.

SOFIA is the premier observatory to study dust in the wavelength regime of its maximum emission.

- How and where is dust formed in explosive events (novae/SNe)?
- How does dust survive in the initial/reverse shocks?
- How is dust formed in outflows of evolved stars?
- How does dust evolve in star-forming dense molecular cores? In proto-planetary disks?









What do we mean by “Interstellar Medium”?

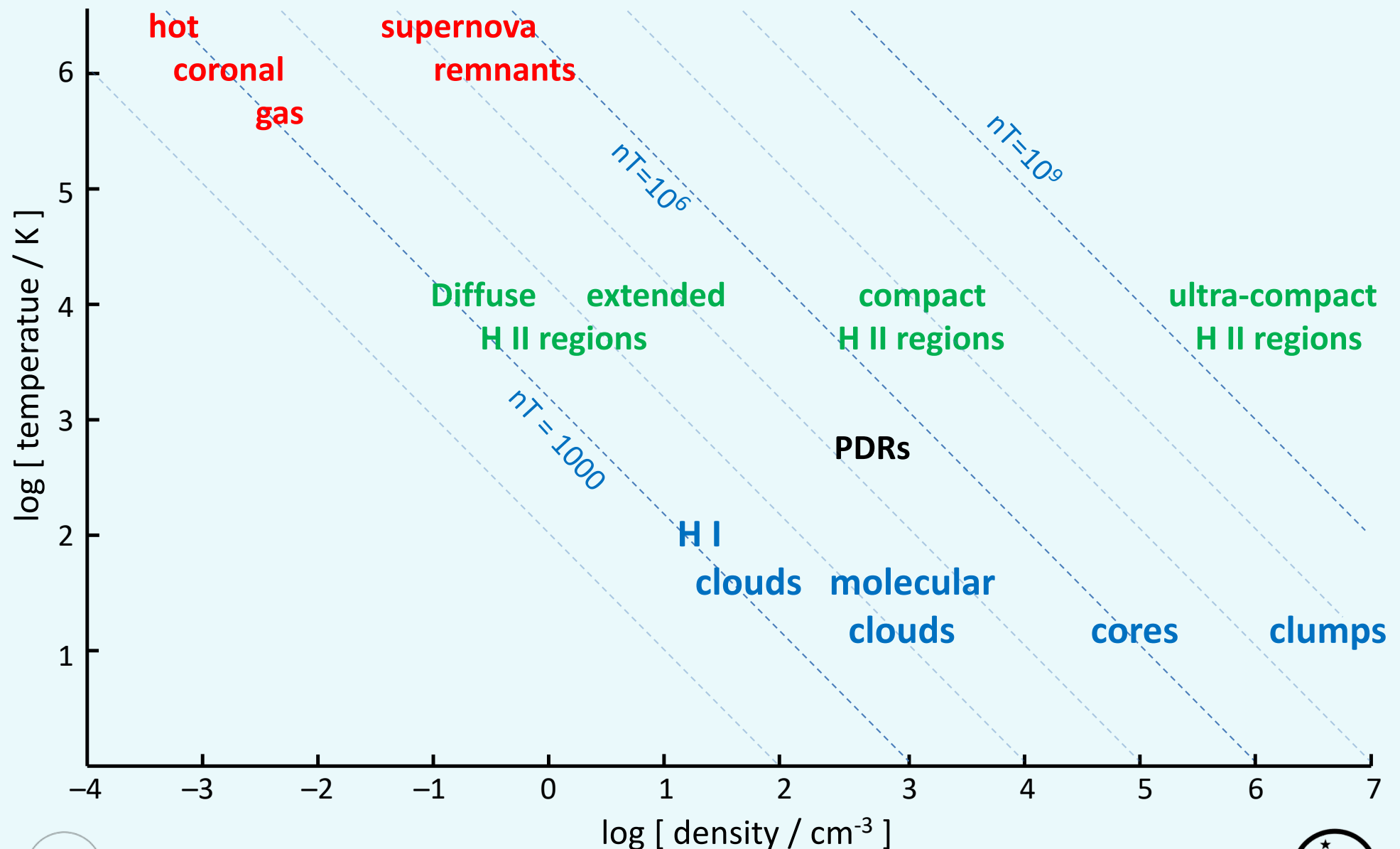


“Everything” in a galaxy except stars, planets, other large bodies and Dark Matter. For the MWG...

- **Gas:** $n \sim 10^{-4} - 10^7 \text{ cm}^{-3}$, $T \sim 5 - 10^7 \text{ K}$
 $\langle nkT \rangle \sim \text{several } 10^{-13} \text{ erg cm}^{-3}$
- **Dust:** $\rho_d \approx 0.01 \rho_{\text{gas}}$, $T \sim 5 - 2000 \text{ K}$
- **Cosmic Rays:** $\zeta \sim 10^{-16} \text{ s}^{-1}$ (several $10^{-13} \text{ erg cm}^{-3}$)
- **Magnetic Fields:** $B \sim 3 \mu\text{G}$ ($4 \times 10^{-13} \text{ erg cm}^{-3}$)
- **Interstellar Radiation Field:** $\sim 10^{-12} \text{ erg cm}^{-3}$



Phases of the Interstellar Medium

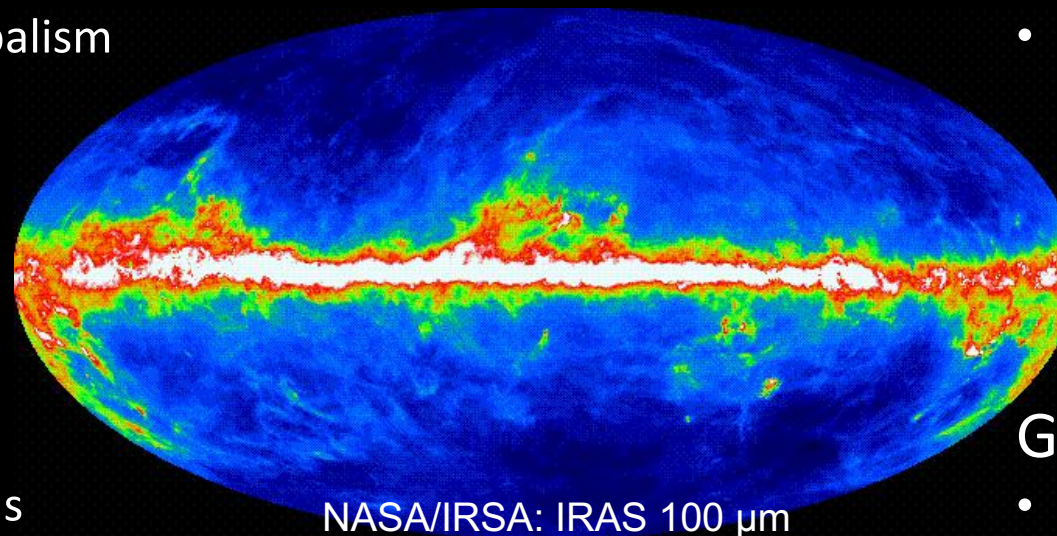


Dust Sources:

- “Superwinds” of AGBs
- SN shocks
- Novae
- Infall from IGM
- Galactic cannibalism

Gas Sources:

- Stellar winds
- SN explosions
- Novae
- Infall from IGM
- Galactic cannibalism

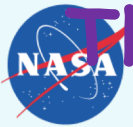


Dust Sinks:

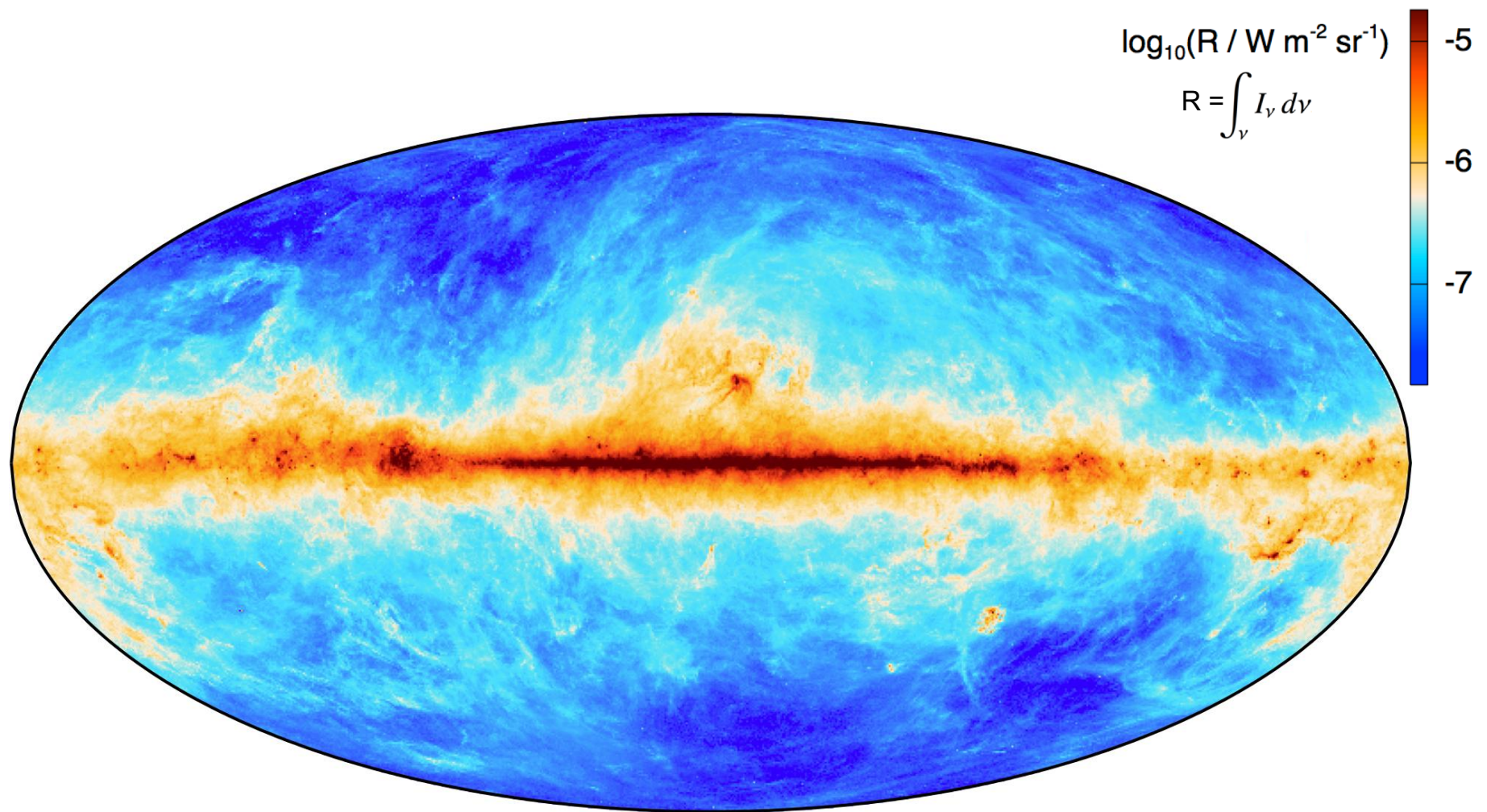
- Dust follows gas
- Destruction in shocks
- Destruction by radiation or CR
- Radiative levitation

Gas Sinks:

- Low mass objects
- Stellar remnants
- Accretion onto stars & central SMBH
- Galactic wind / jet



The Big Picture: Distribution of Dust Emission



arXiv:1312.1300v2

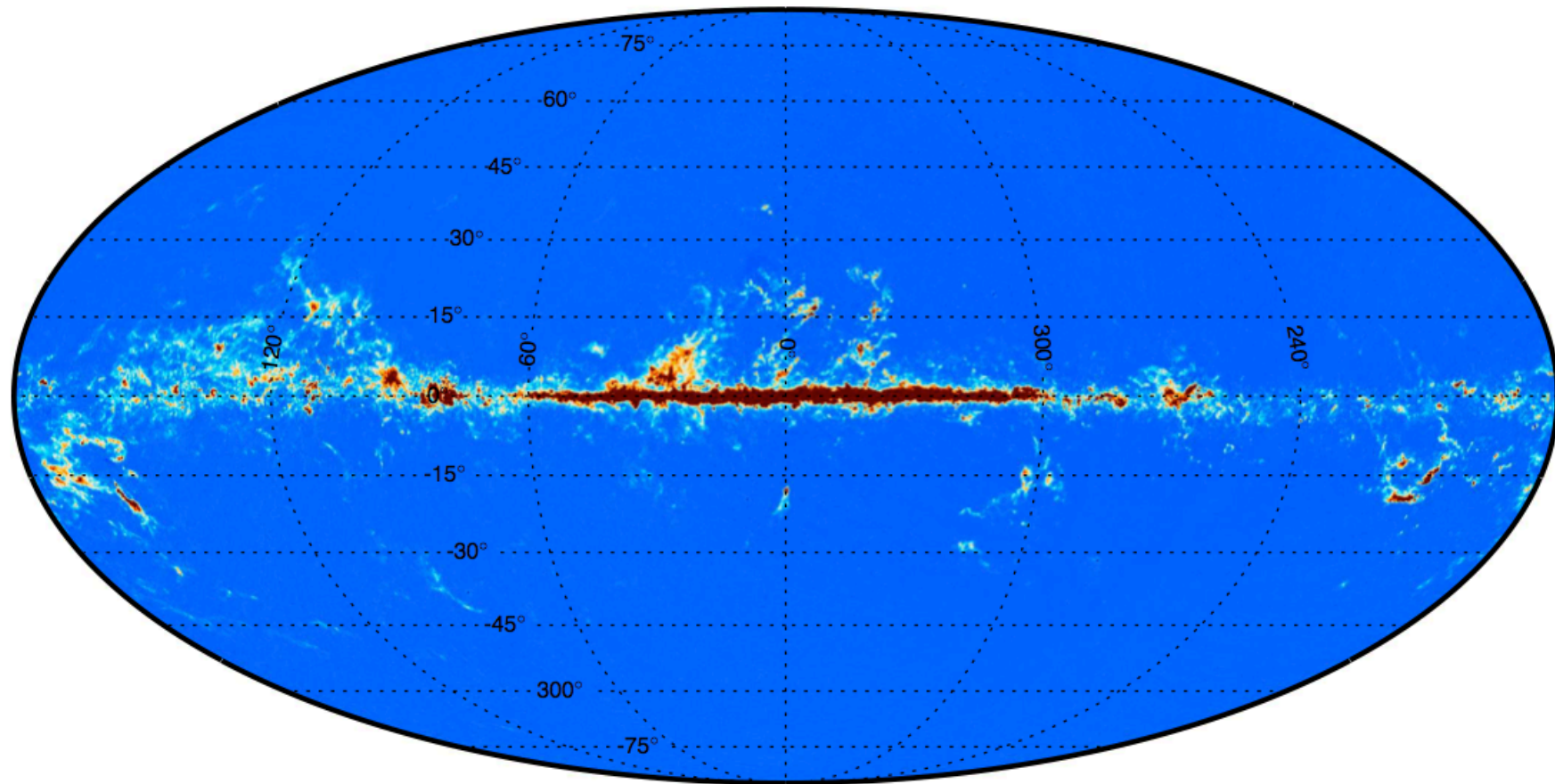
Image credit: ESA/Planck Collaboration



The Big Picture: Distribution of CO emission



Type 3 J=1-0 CO map



-5.0 30.0

K Km/s

arXiv:1303.5073v2

Image credit: ESA/Planck Collaboration

Basic Equation: $d/ds(I_V) = -\sigma_V (I_V - S_V)$

Radiation Intensity: I_V

Source function: S_V

detailed physics

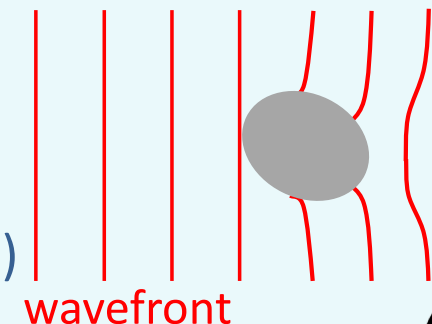
Extinction coefficient: σ_V Often use (gram-)opacity: $\kappa_V = \sigma_V/\rho$

$$\sigma_V = \sigma_V^{sca} + \sigma_V^{abs}$$

- Dust often treated “classically” (this does not work for PAHs)

$$\sigma_V S_V = \sum_i \sigma_{V,i}^{abs} B_V(T_i) \quad (\text{sum over lots of “grey bodies”})$$

- Grain parameters that affect $\sigma_{V,i}$ (e.g. Mie theory of scattering):
 - grain composition: complex dielectric constants
=> index of refraction: $m = n + ik$
 - grain size a_i (can be normalized to λ/a)
 - grain shape and orientation (polarization properties)



wavefront



Abundance of elements & ionization potentials



| | | Log Abundance | | Ionization Potentials | | |
|----------|-----------|---------------|-------|-----------------------|--------|--------|
| Element | Atomic Wt | number | mass | I | II | III |
| 1 H | 1.008 | 12.00 | 12.00 | 13.598 | | |
| 2 He | 4.003 | 10.99 | 11.59 | 24.587 | 54.416 | |
| 6 C | 12.011 | 8.56 | 9.64 | 11.260 | 24.383 | 47.887 |
| 7 N | 14.007 | 8.05 | 9.19 | 14.534 | 29.601 | 47.448 |
| 8 O | 15.999 | 8.72 | 9.92 | 13.618 | 35.117 | 54.934 |
| 10 Ne | 20.180 | 8.09 | 9.34 | 21.564 | 40.962 | 63.45 |
| 11 Na | 22.990 | 6.31 | 7.67 | 5.139 | 47.286 | 71.64 |
| 12 Mg | 24.305 | 7.59 | 8.97 | 7.646 | 15.035 | 80.143 |
| 13 Al | 26.982 | 6.48 | 7.91 | 5.986 | 18.826 | 28.228 |
| 14 Si | 28.086 | 7.56 | 9.00 | 8.151 | 16.345 | 33.492 |
| 16 S | 32.065 | 7.25 | 8.75 | 10.360 | 23.33 | 34.83 |
| 18 Ar | 39.948 | 6.8 | 8.4 | 15.759 | 27.629 | 40.74 |
| 20 Ca | 40.078 | 6.30 | 7.93 | 6.113 | 11.871 | 50.91 |
| 26 Fe | 55.847 | 7.60 | 9.35 | 7.870 | 16.16 | 30.651 |
| 28 Ni | 58.693 | 6.3 | 8.04 | 7.635 | 18.168 | 35.17 |
| Σ (rest) | <55.08> | 6.26 | 8.00 | | | 14 |



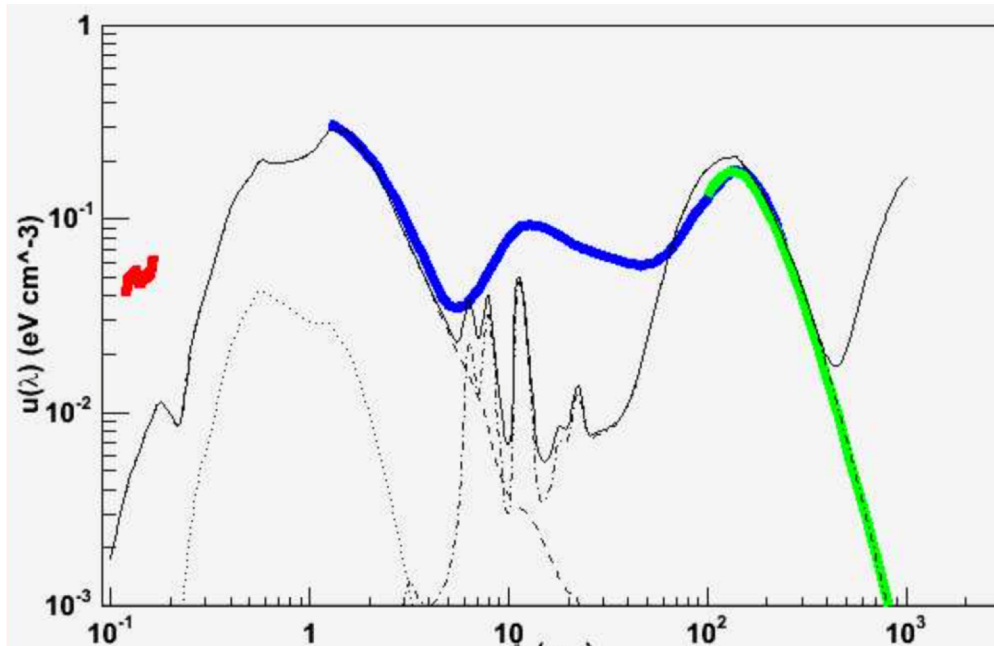
What do we learn from this table?



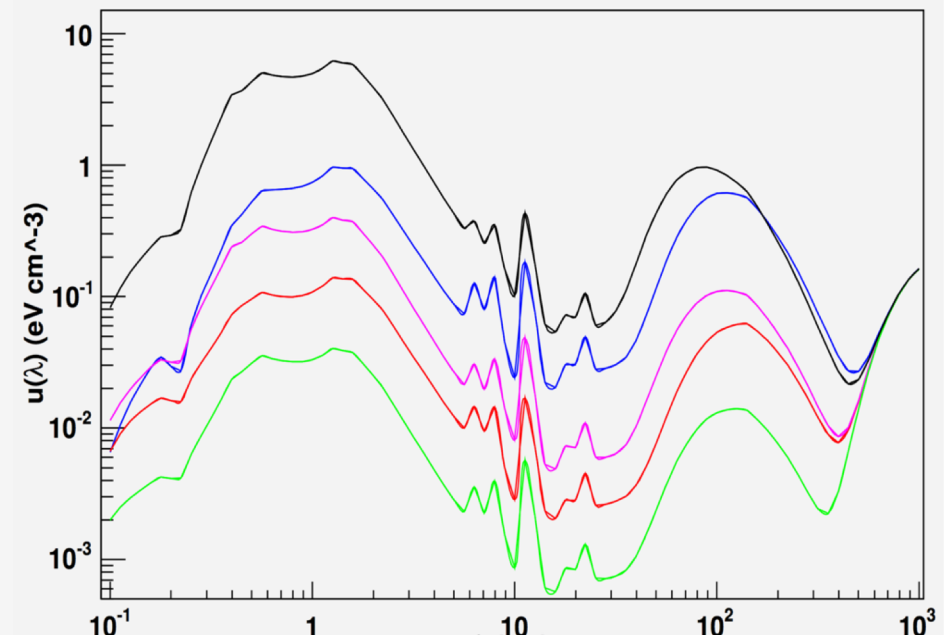
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- O is neutral whenever H is neutral
- C, Na, Mg, Al, Si, S, Ca, Fe, Ni can all be easily ionized by background starlight (Ca doubly), providing a floor for the relative electron abundance
- Among the most 6 abundant elements ($\log N_x / \log N_H > -4$), only C can be (singly) ionized in regions where H is neutral
- OIII traces very energetic UV fields
- Abundances provide constraints on dust composition





Wavelength [μm]



Wavelength [μm]

Local radiation field as modelled by Porter and Strong (2005)
Solid black line total radiation field, including CMBR. *Dashed black line*: total optical. *Dotted black line*: total scattered optical. *Dash-dot black line*: total infra-red.
 Data: *thick red solid line*: Apollo; *thick blue solid line*: DIRBE; *thick green solid line*: FIRAS.

Spatial variation of the total radiation field as a function of galactocentric distance.

black: $R = 0$ kpc, $z = 0$
 blue: $R = 5$ kpc, $z = 0$
 magenta: $R = 0$ kpc, $z = 5$ kpc
 red: $R = 12$ kpc, $z = 0$
 green: $R = 20$ kpc, $z = 0$



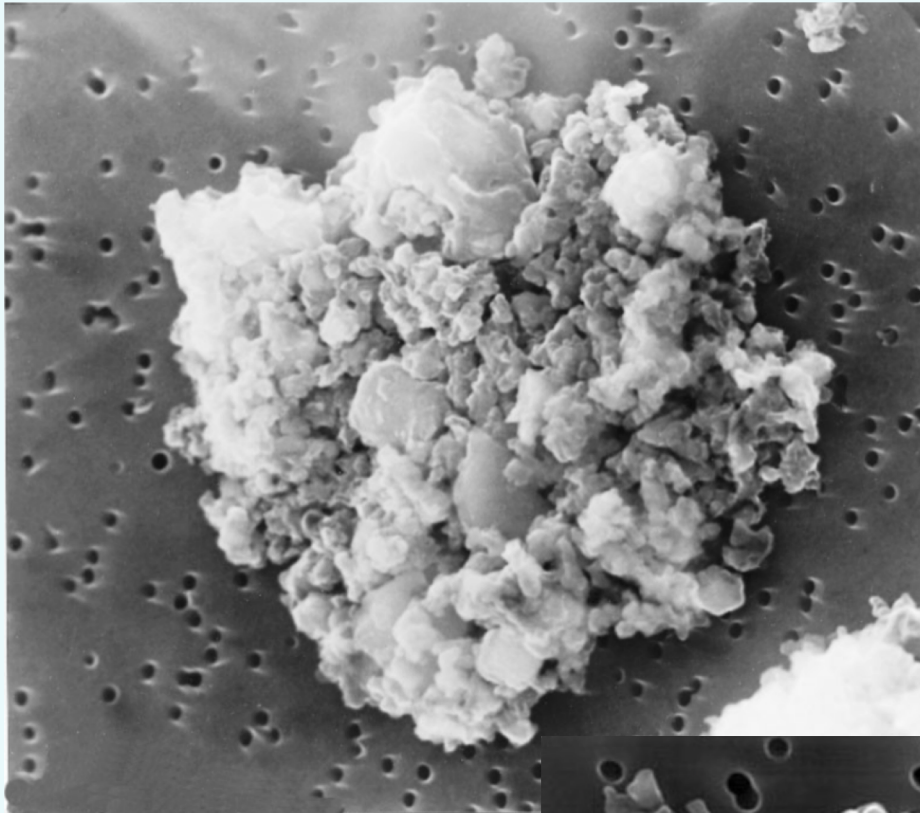
What do we want to know about interstellar grains?



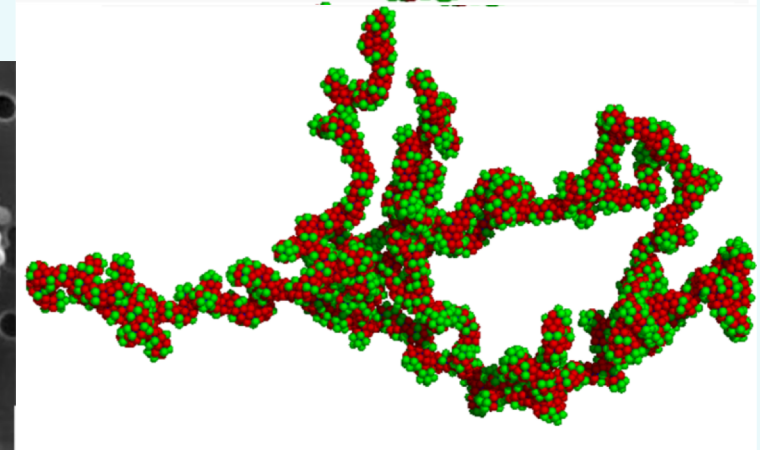
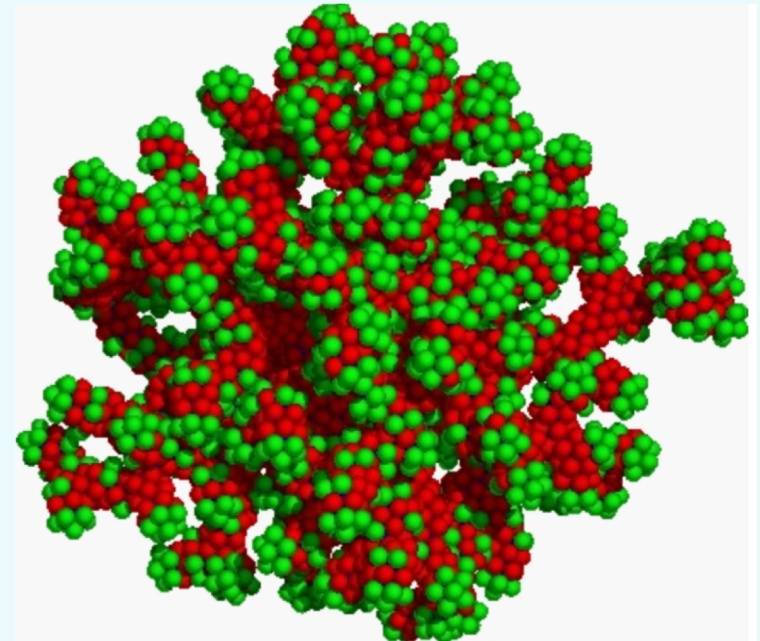
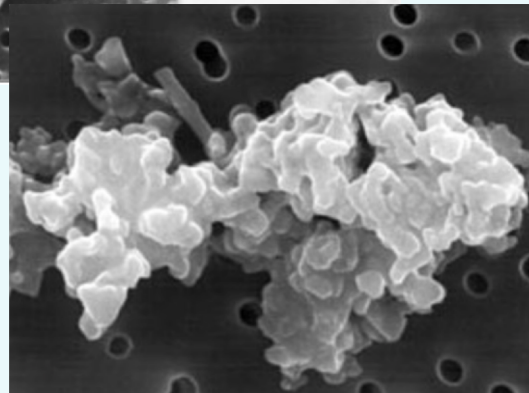
- Their composition
 - Consistent with observed abundances (comparison between stellar and interstellar abundances: depletion)
 - Consistent with the solid-state physics (condensation sequence, thermodynamic stability of the components, structural changes)
 - Constrained by solid state absorption/emission features (specific to the solid, its structure)
 - Should be the precursors of pre-solar (meteoritic) grains
- Their size distribution
 - Can be constrained from: Extinction
 - Emission (shape of the emission features)
- How they interact with light
 - Scattering properties (polarization, when aligned)
 - Absorption (and emission) properties
- Their role in the thermodynamics (heating/cooling) of the gas
- Their role in surface chemistry and adsorption of species
- Their role in the hydrodynamics of the gas



(Un)fortunately, grains are not spheres



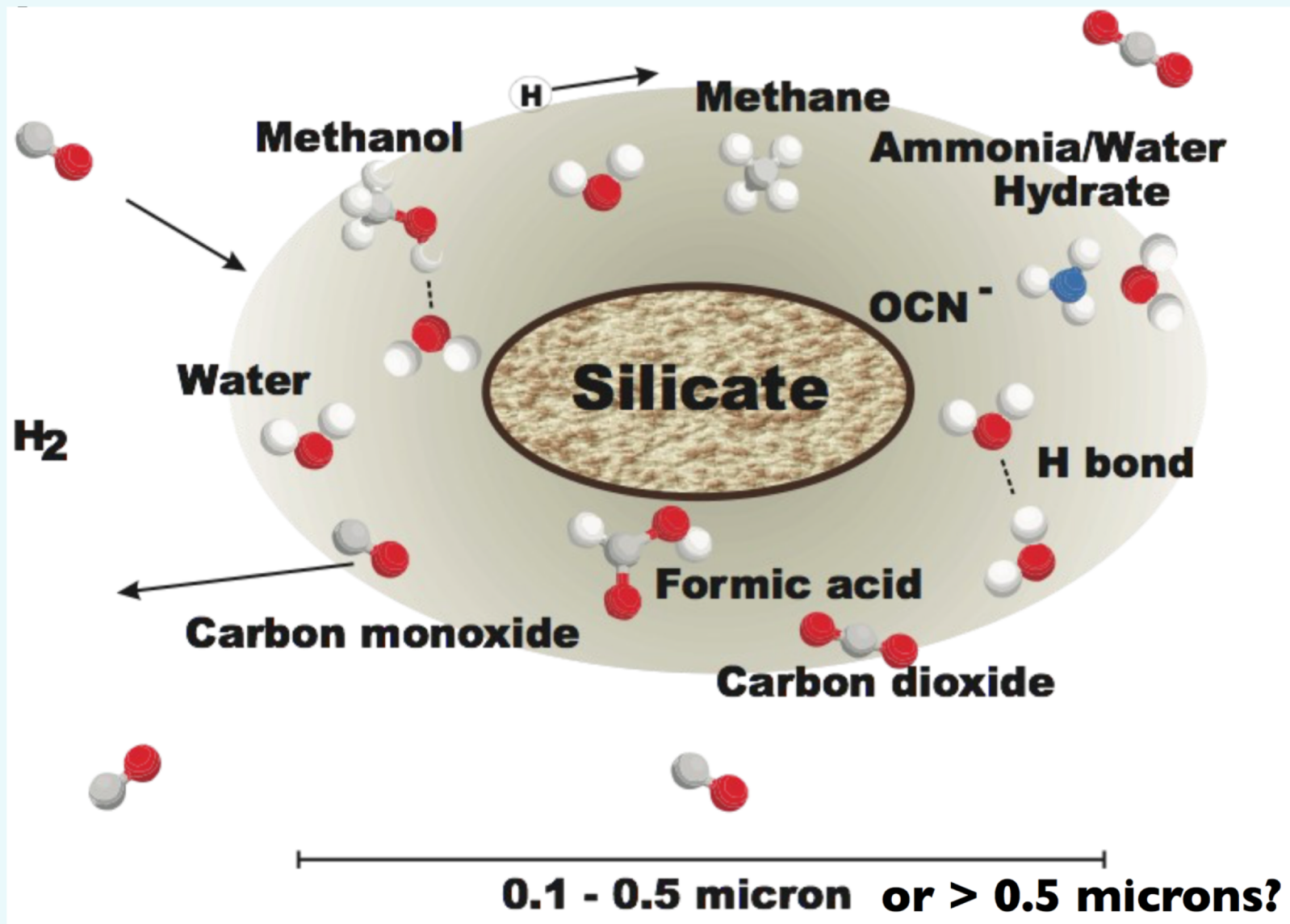
IDP Image Credit NASA



- There cannot be much more dust than $0.01 M_{\text{H}}$, because that uses up all elements
- Typical model:
 - 2/3 of carbon used for carbonaceous material
 - Essentially all Mg, Fe, Si and 20% of O in $(\text{Mg,Fe})_2\text{SiO}_4$ consistent with the observational depletion pattern and the physics of solid condensation
 - Some SiC if grains are formed in carbon-rich environments
 - May include mantel of “dirty ices” (volatiles)



Icy dust models (mantel mostly H_2O)



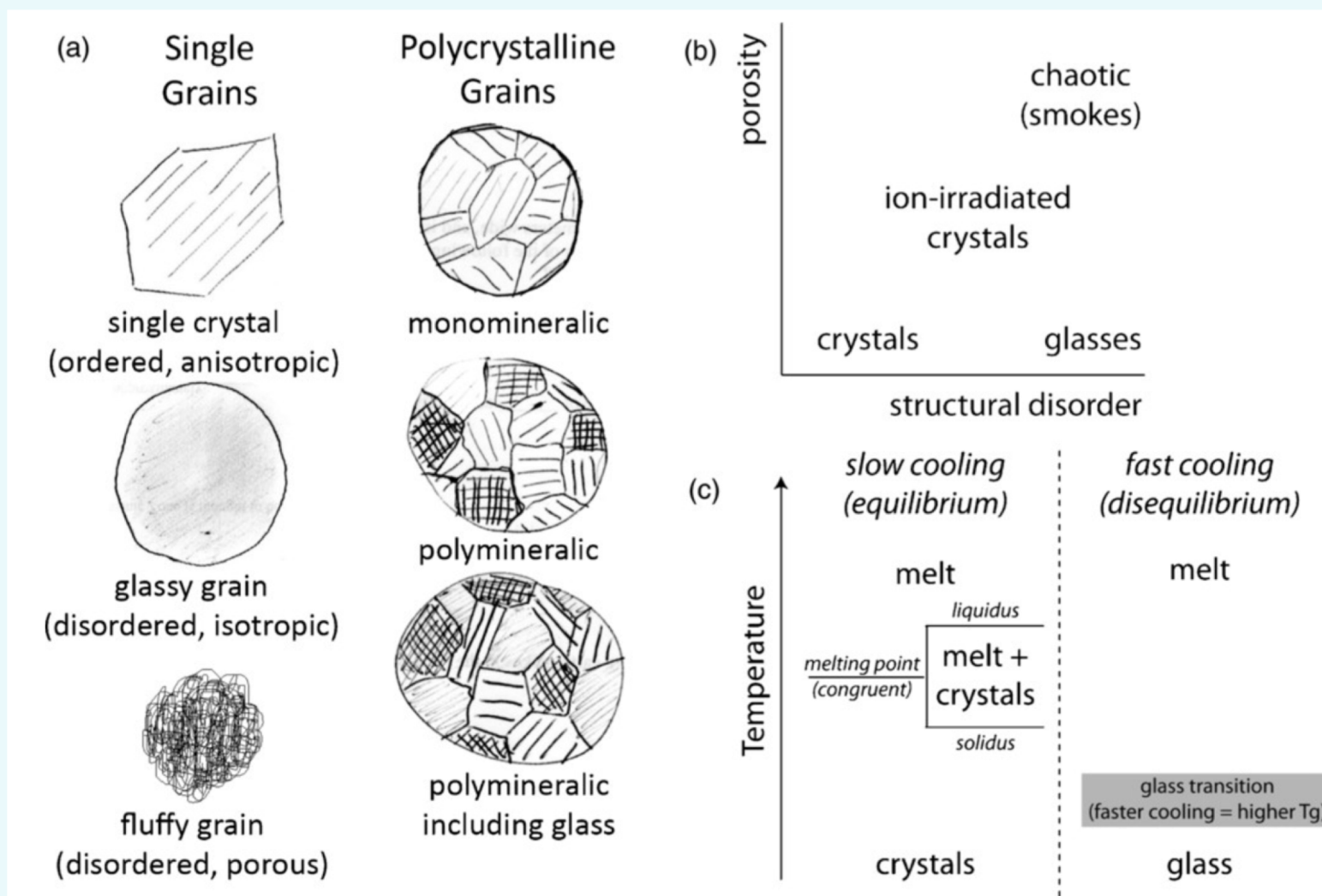
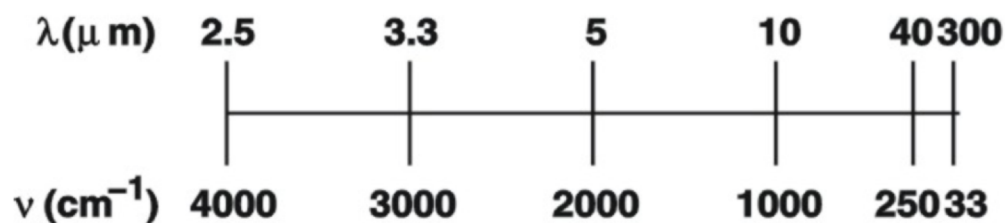
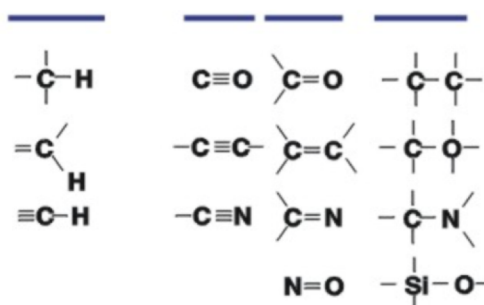


Figure 1 from Disordered Silicates in Space: A Study of Laboratory Spectra of "Amorphous" Silicates Angela K. Speck et al. 2011 ApJ 740 93

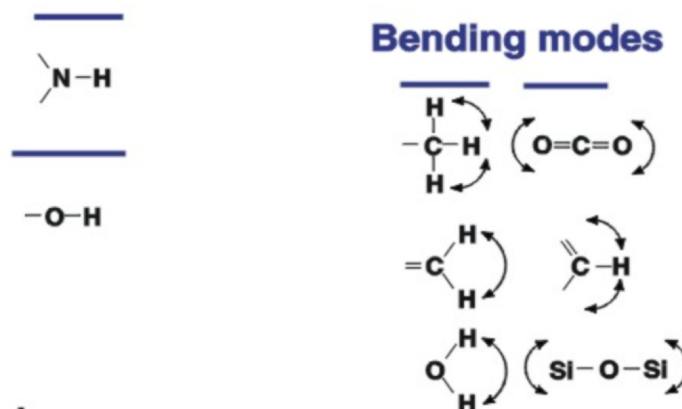


← Stretching modes →

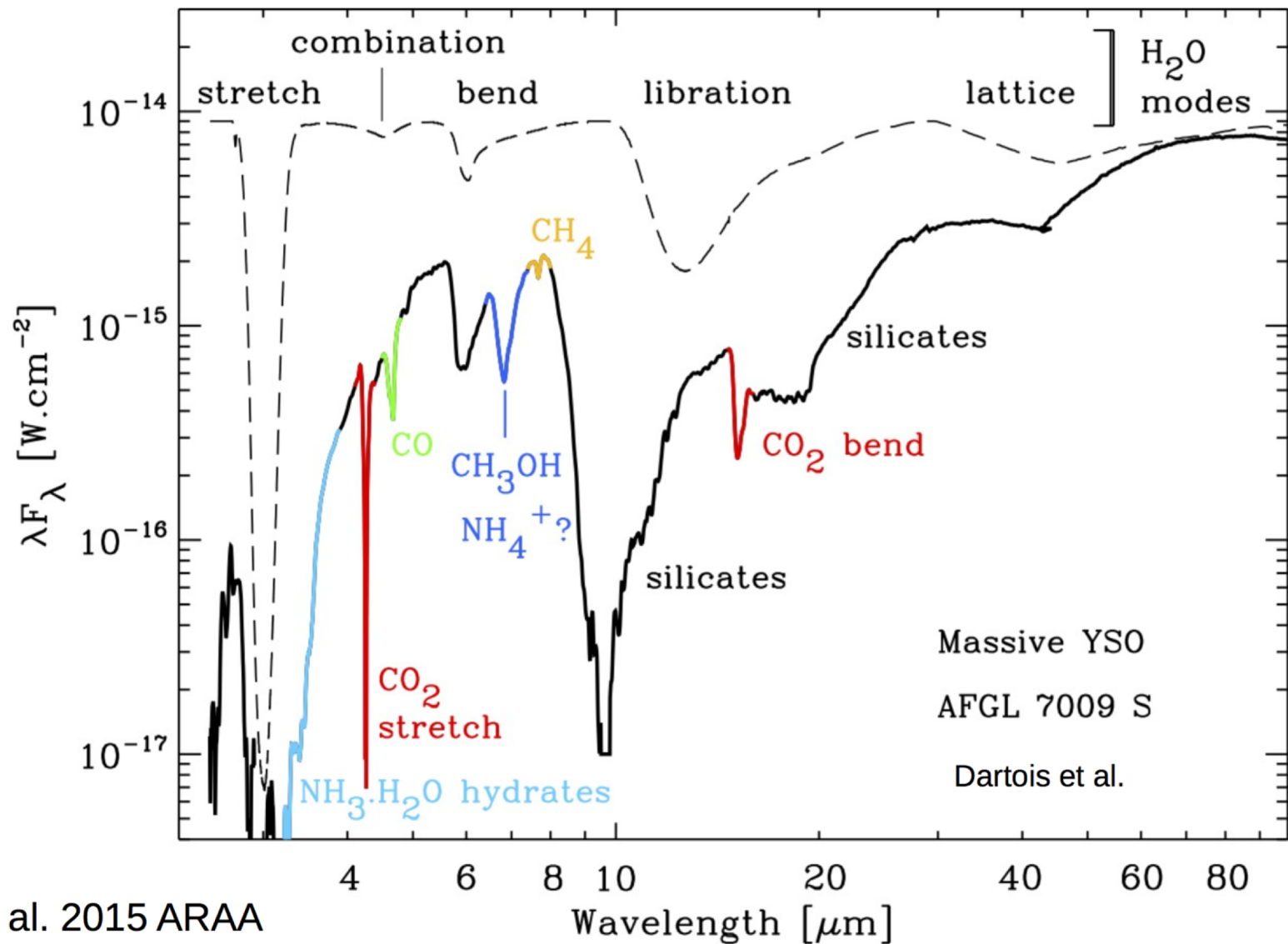
Lattice modes



Bending modes



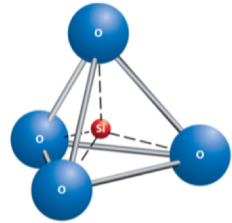
- Stretching & bending modes occur in both amorphous and crystalline solids
- Lattice (phonon) modes exist only in crystalline solids



Boogert et al. 2015 ARAA



Amorphous vs Crystalline silicate: structure and spectral signatures

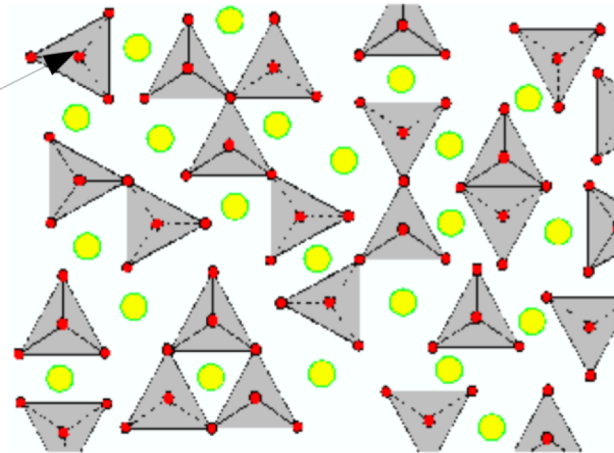


Henning T. 2010.
Annu. Rev. Astron. Astrophys. 48:21–46

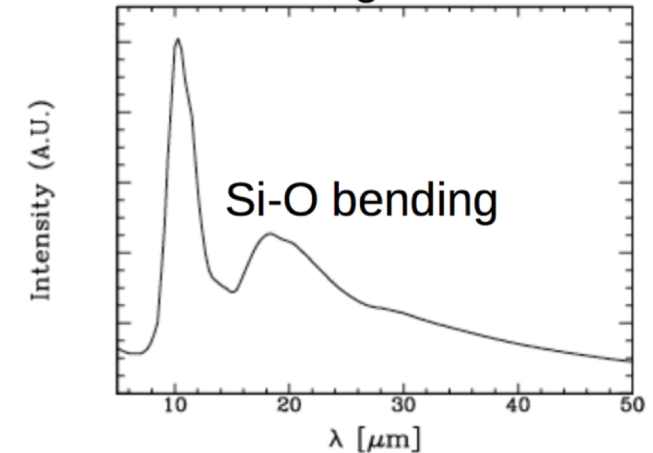


Molster et al. 2005

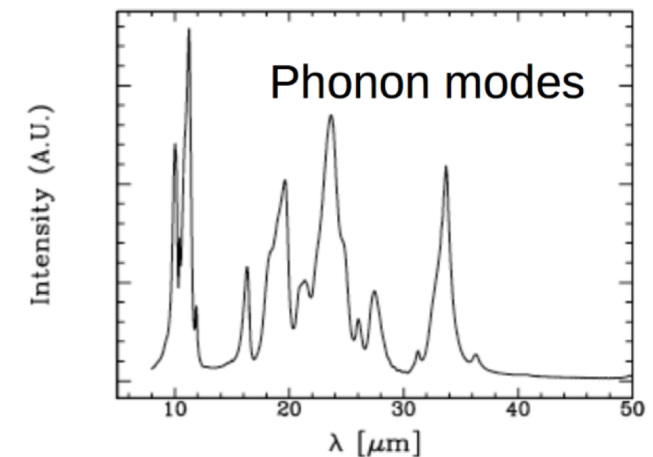
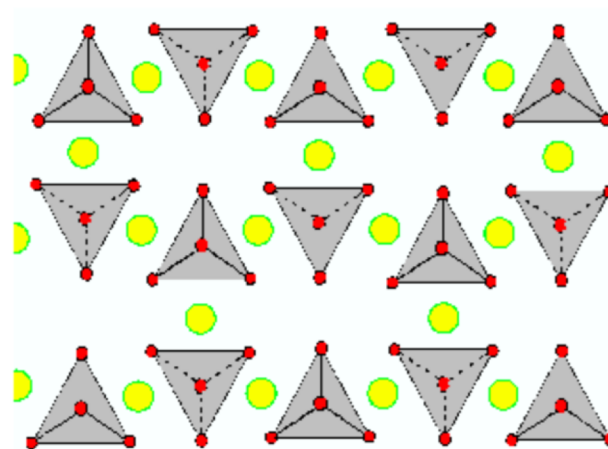
Amorphous structure

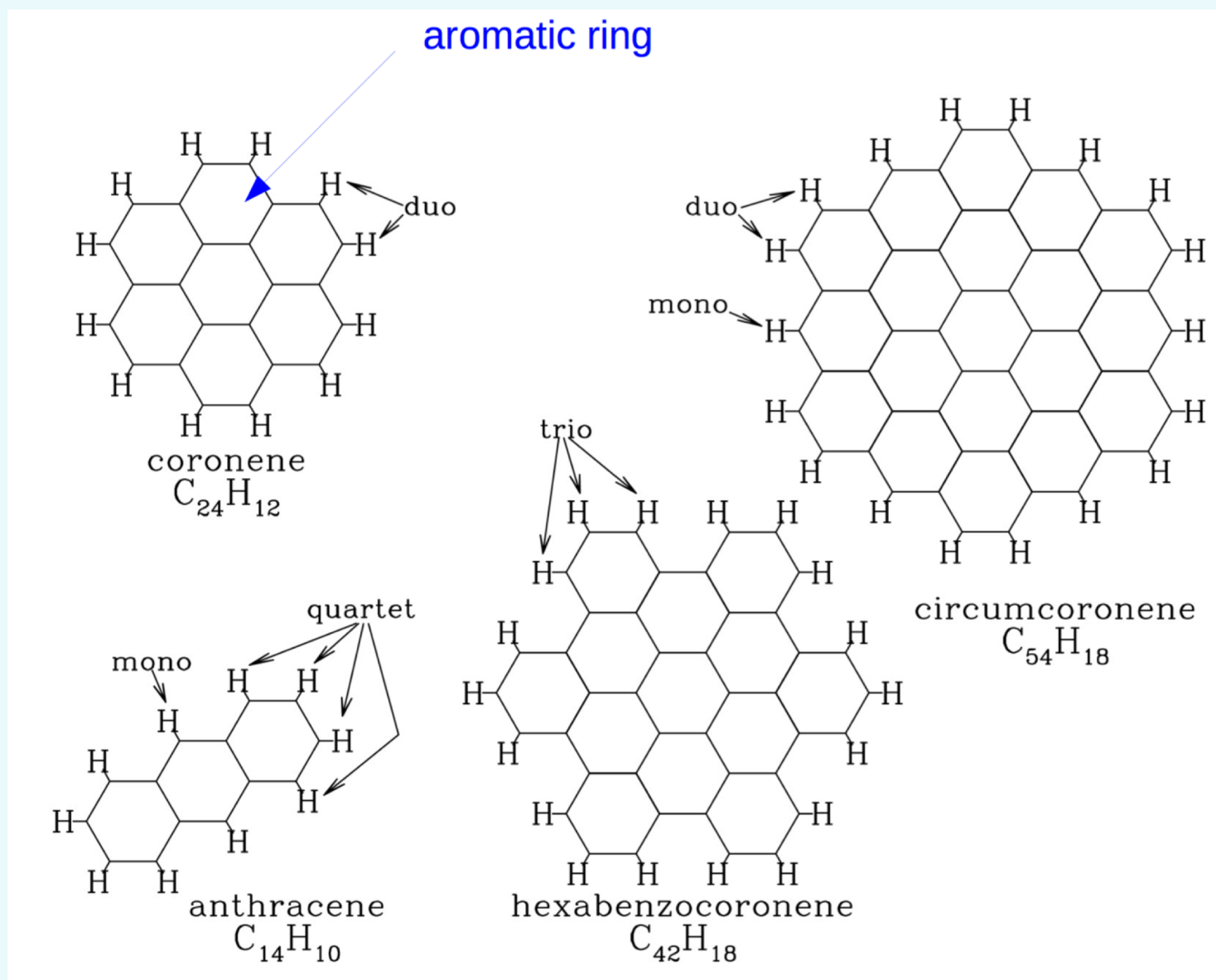


Si-O stretching



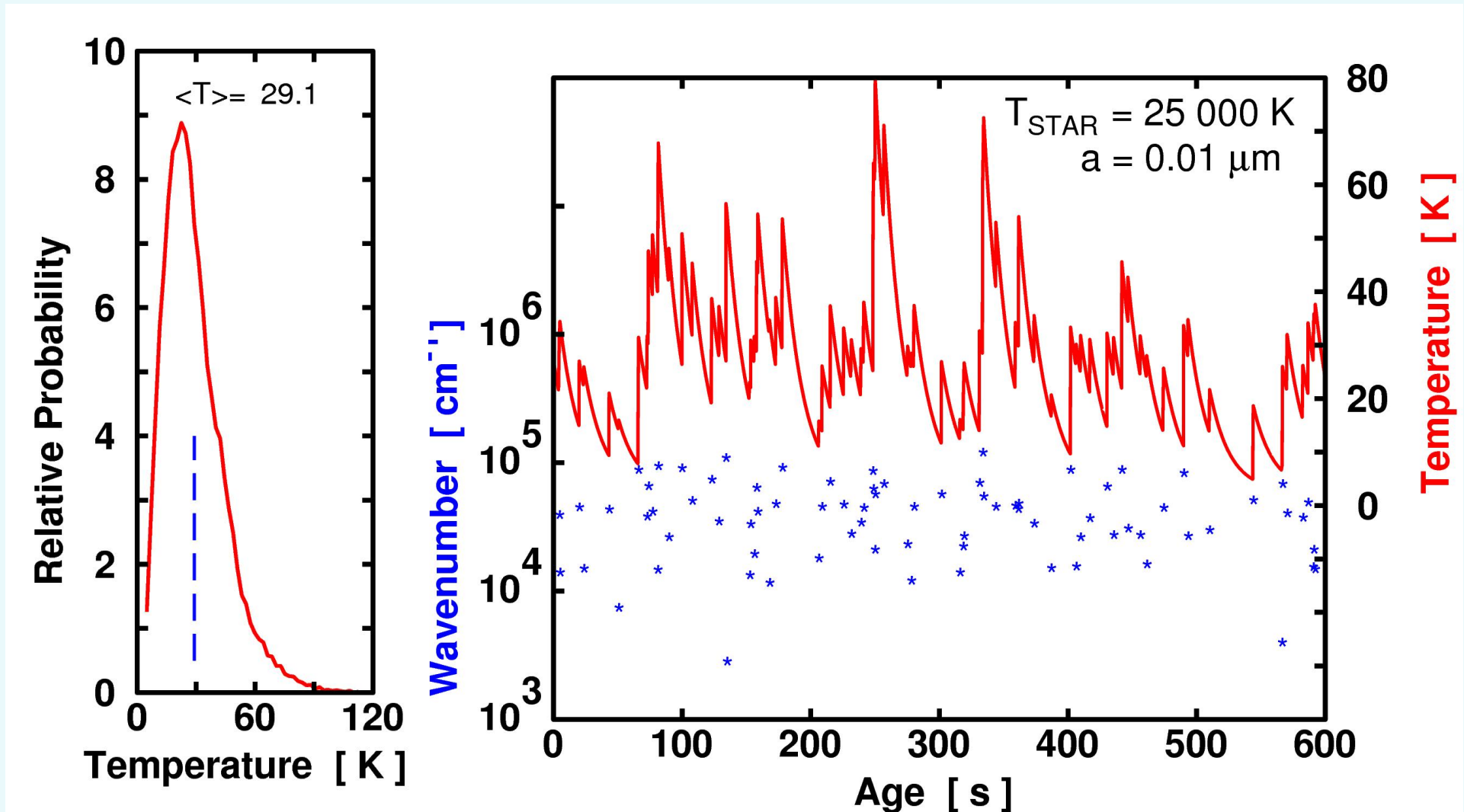
Crystalline structure







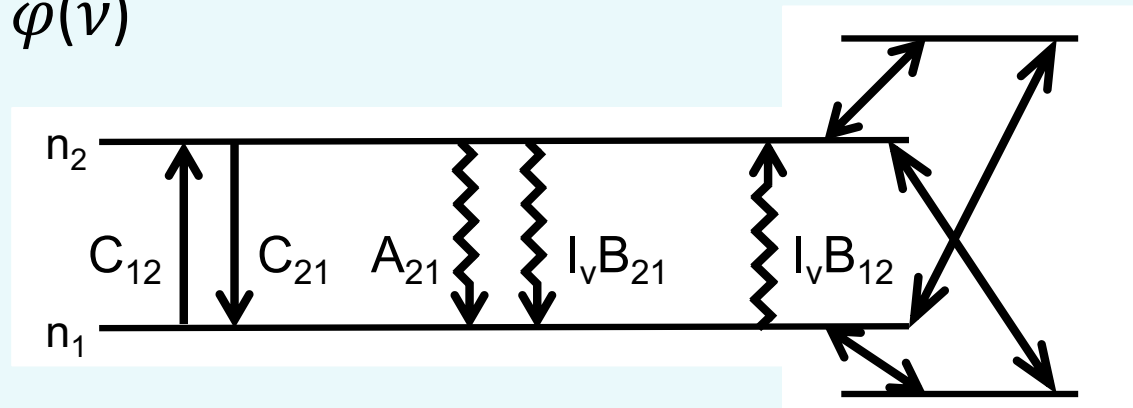
Stochastic heating of small grains



Basic Equation:
$$d/ds(I_V) = - \underbrace{\sigma_V^C}_{\text{detailed physics}} (I_V - \underbrace{S_V^C}_{\text{more detailed physics}}) - \underbrace{\sigma_V^L}_{\text{detailed physics}} (I_V - \underbrace{S_V^L}_{\text{more detailed physics}})$$

$$\sigma_V^L = n_1 B_{12} h\nu_0/4\pi [1 - n_2 g_1/n_1 g_2] \varphi(\nu)$$

$$\sigma_V^L S_V^L = n_2 A_{21} h\nu_0/4\pi \varphi(\nu)$$



A_{21} , B_{21} , B_{12} are Einstein coefficients

$\varphi(\nu)$ is the line profile (includes thermal doppler broadening)



The multi-level case



$$n_i \left(\sum_{j<i} A_{ij} + \sum_{j\neq i} B_{ij} u_{ij} + \sum_{j\neq i} C_{ij} \right) = \sum_{j>i} A_{ji} n_j + \sum_{j\neq i} B_{ji} u_{ij} n_j + \sum_{j\neq i} C_{ij} n_j$$

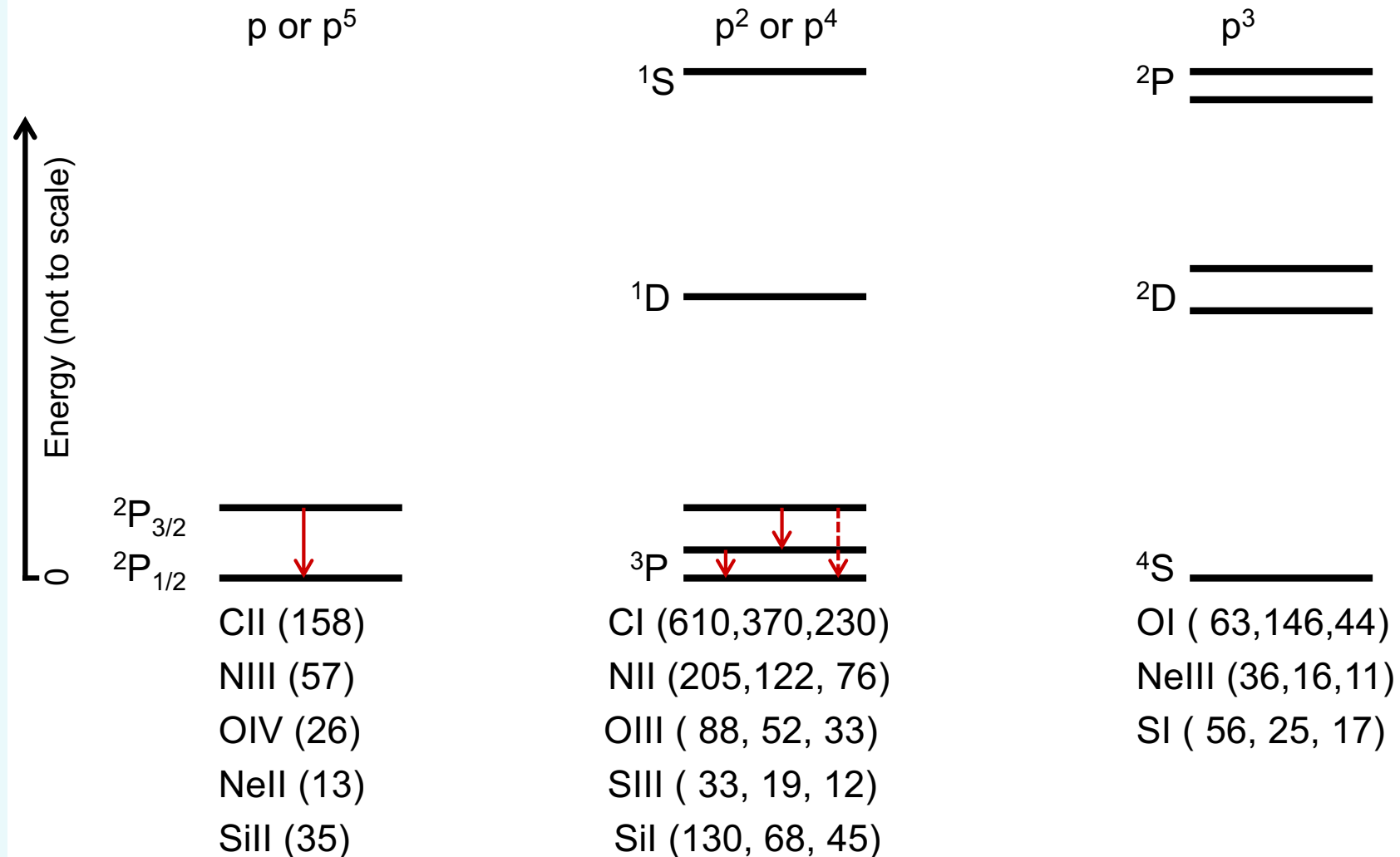
where $u_{ij} = \int_{4\pi} I(\nu_{ij}) d\Omega$ is the radiation density
and $\sum_i n_i$ = density of species



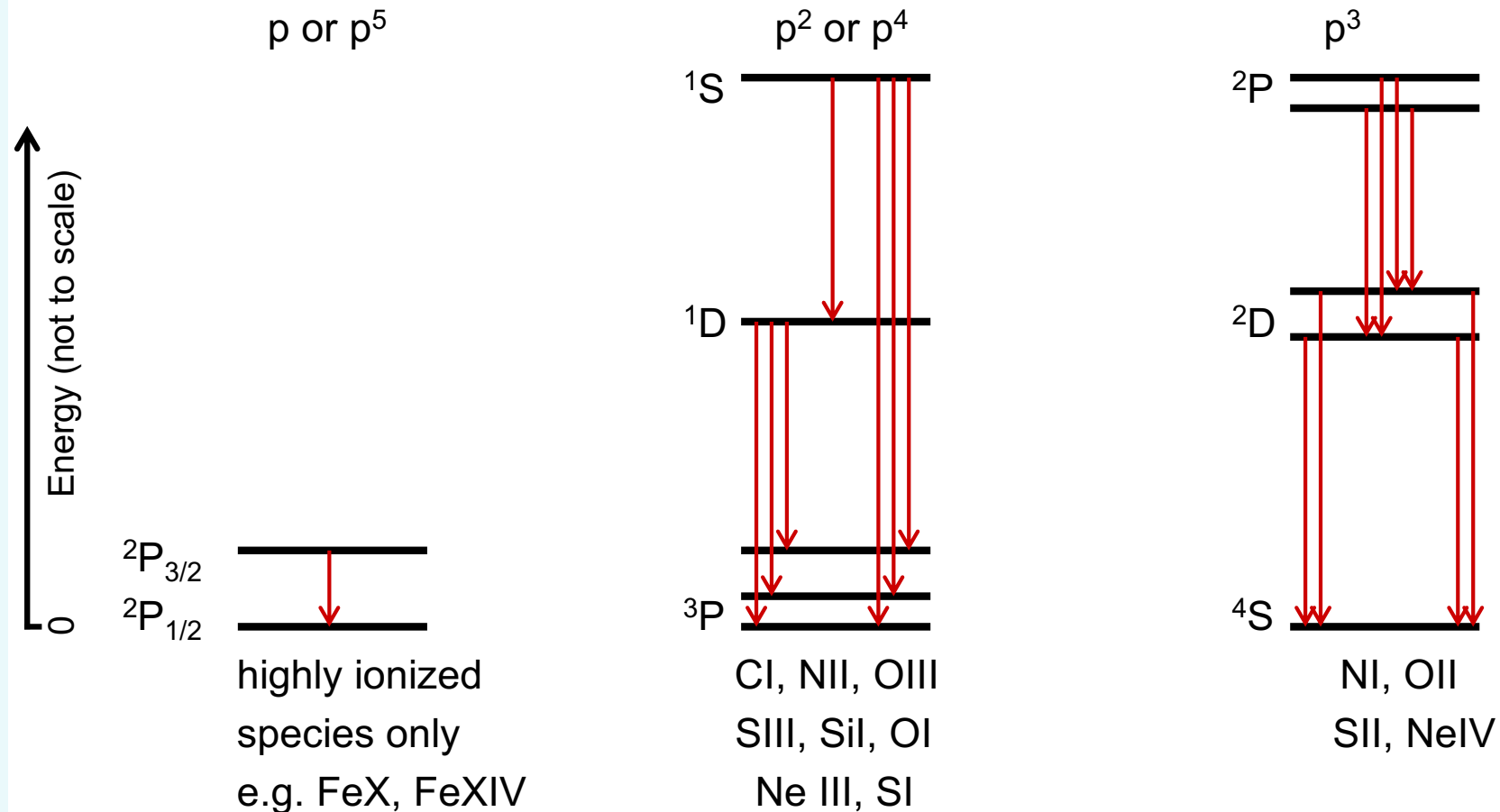
Fine Structure Line Emission



L-S coupling for ions with one to five p electrons: (far)IR transitions

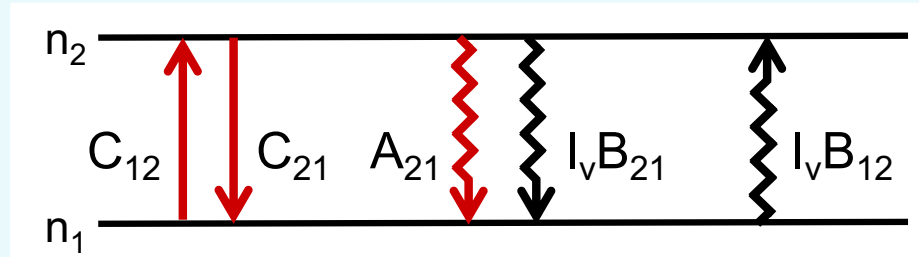


L-S coupling for ions with one to five p electrons: transitions in visible



$$C_{12}/C_{21} = g_2/g_1 \exp(-\Delta E/kT)$$

$$n_1 C_{12} = n_2 (C_{21} + \beta_{21} A_{21})$$



$n = n_1 + n_2$, where n is the density of coolant available

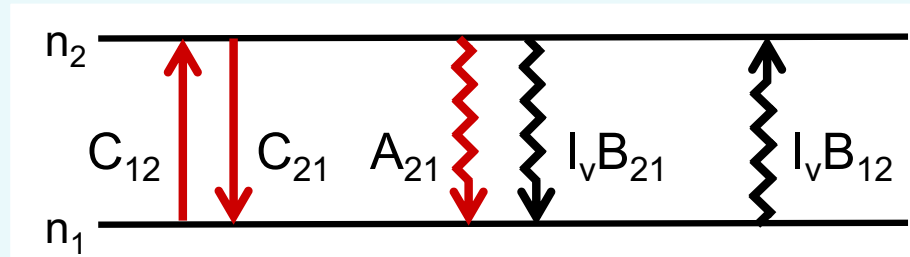
The cooling rate from this transition is: $\Lambda_{21} = \Delta E n_2 A_{21} \beta_{21}$

β_{21} is the “escape probability”, i.e. the likelihood that the emitted photon can actually leave the volume without being reabsorbed. Often one uses the simple expression $\beta_{21} \sim \langle [1 - \exp(-\tau_{21})] / \tau_{21} \rangle$, where τ_{21} is the line center optical depth

$$\Lambda_{21} = \Delta E n A_{21} \beta_{21} \frac{g_2/g_1 \exp(-T_{\text{ex}}/T)}{1 + g_2/g_1 \exp(-T_{\text{ex}}/T) + n_{\text{ncrit}}/n_{\text{coll}}} ,$$

where $n_{\text{crit}} = n_{\text{coll}} A_{21} \beta_{21} / C_{21} = A_{21} \beta_{21} / \langle v\sigma \rangle_{21}$ is the “critical density” and n_{coll} is the density of the exciting collider

$$n_{\text{coll}} \ll n_{\text{crit}} \text{ and } n_{\text{coll}} \gg n_{\text{crit}}$$



$$\Lambda_{21} = \Delta E n A_{21} \beta_{21} \frac{g_2/g_1 \exp(-T_{\text{ex}}/T)}{1 + g_2/g_1 \exp(-T_{\text{ex}}/T) + n_{\text{ncrit}}/n_{\text{coll}}}$$

$$n_{\text{coll}} \ll n_{\text{crit}} \text{ (subthermal)}$$

$$n_{\text{coll}} \gg n_{\text{crit}} \text{ (thermalized)}$$

$$\Lambda_{21} = \Delta E n n_{\text{coll}} \langle v\sigma \rangle_{21}$$

$$\Lambda_{21} = \Delta E n A_{21} \beta_{21} \frac{g_2/g_1 \exp(-T_{\text{ex}}/T)}{1 + g_2/g_1 \exp(-T_{\text{ex}}/T)}$$

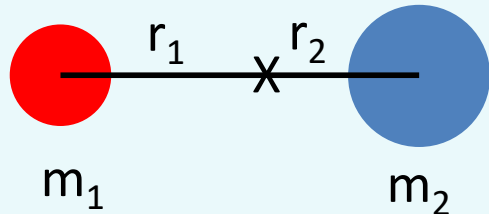
Examples (subthermal):

$$\Lambda_{\text{CII}}(158\mu\text{m}) = (6.67 \times 10^{-20} n_e T^{-1/2} + 1.77 \times 10^{-23} n_{\text{HI}} T^{1/2}) n_{\text{CII}} \exp(-92\text{K}/T)$$

$$\Lambda_{\text{OI}}(63\mu\text{m}) = 6.43 \times 10^{-24} n_{\text{HI}} n_{\text{OI}} T^{1/2} \exp(-228\text{K}/T)$$

$$\Lambda_{\text{OI}}(147\mu\text{m}) = 4.7 \times 10^{-24} n_{\text{HI}} n_{\text{OI}} T^{1/2} \exp(-326\text{K}/T)$$

Assume two nuclei, m_1 & m_2 separated by distance d



$$r_1 + r_2 = d$$

$$r_1 = d m_2 / (m_1 + m_2)$$

$$m_1 r_1 = m_2 r_2$$

$$r_2 = d m_1 / (m_1 + m_2)$$

$$E = \frac{1}{2} [m_1 (r_1 \omega)^2 + m_2 (r_2 \omega)^2]$$

$$L = m_1 (r_1)^2 \omega + m_2 (r_2)^2 \omega$$

$$E = \frac{1}{2} L \omega$$

$$L = \mu d^2 \omega \quad \text{with } \mu = m_1 m_2 / (m_1 + m_2)$$

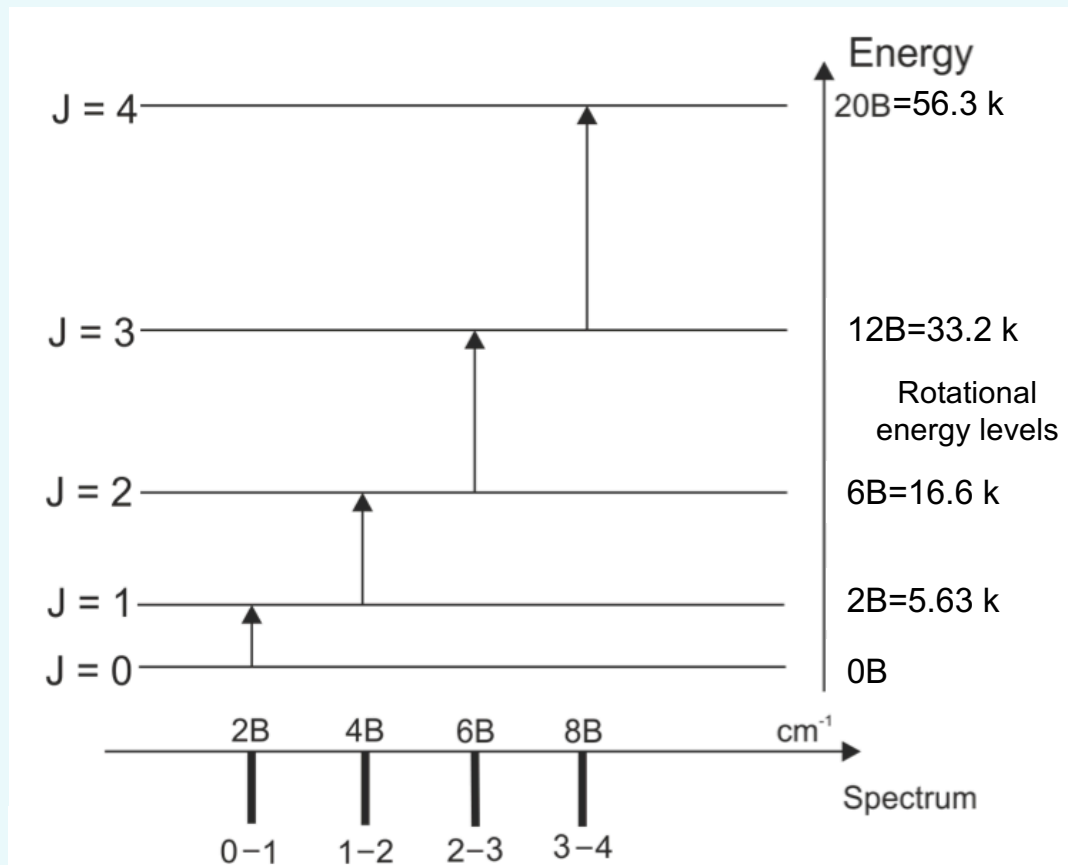
“Quantum” ansatz: $L_n = n \hbar \quad \Rightarrow \quad E_n = n^2 \frac{\hbar^2}{2 \mu d^2}$

Actually, $E_J = J(J+1) \frac{\hbar^2}{2 \mu d^2} = J(J+1) B$

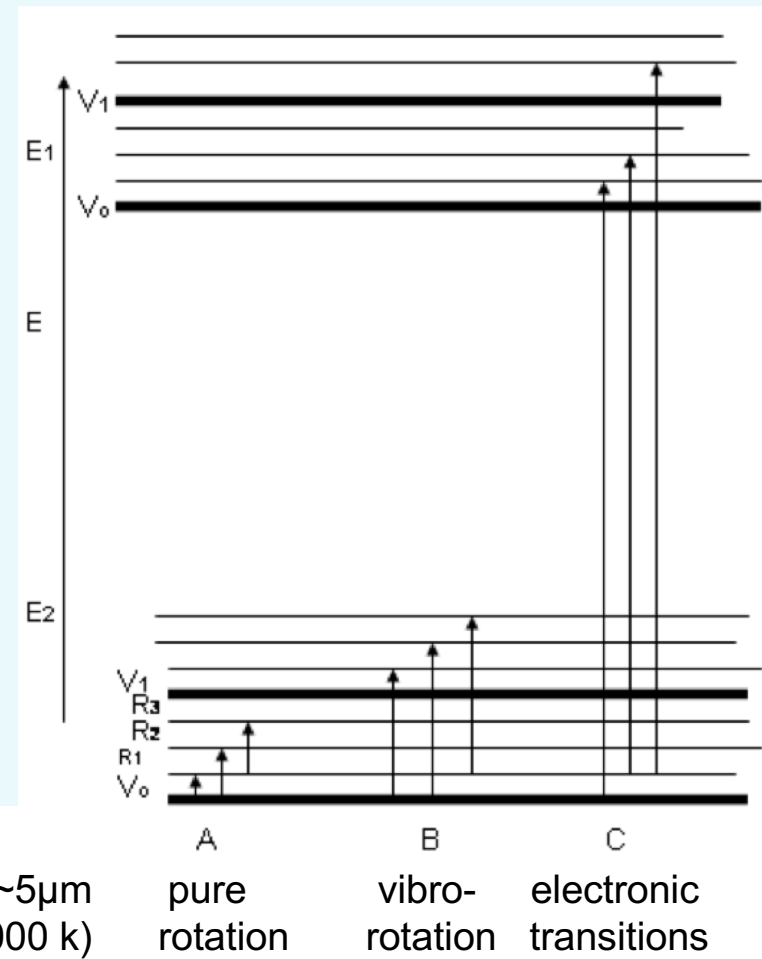
Example: The carbon monoxide molecule (CO) consists of a carbon atom and an oxygen atom separated by a distance of $d = 1.13 \times 10^{-10} \text{ m}$ (from Google).

$$\mu = 6.86 m_H \quad \Rightarrow \quad B = 2.77 k_B = 57.6 \text{ GHz } h = 0.00024 \text{ eV}$$

- CO rotation energy levels $E_J = J(J+1)B$, where $\Delta J = \pm 1$
 - $B \approx \text{const.} = 2.77 \text{ k} = 57.6 \text{ GHz } h = 0.00024 \text{ eV}$ (Actually, B gets larger with J)



The CO $v=1 \rightarrow 0$ rot-vib transitions occur at $\sim 5 \mu\text{m}$ ($T_{\text{ex}} \approx 3000 \text{ K}$)



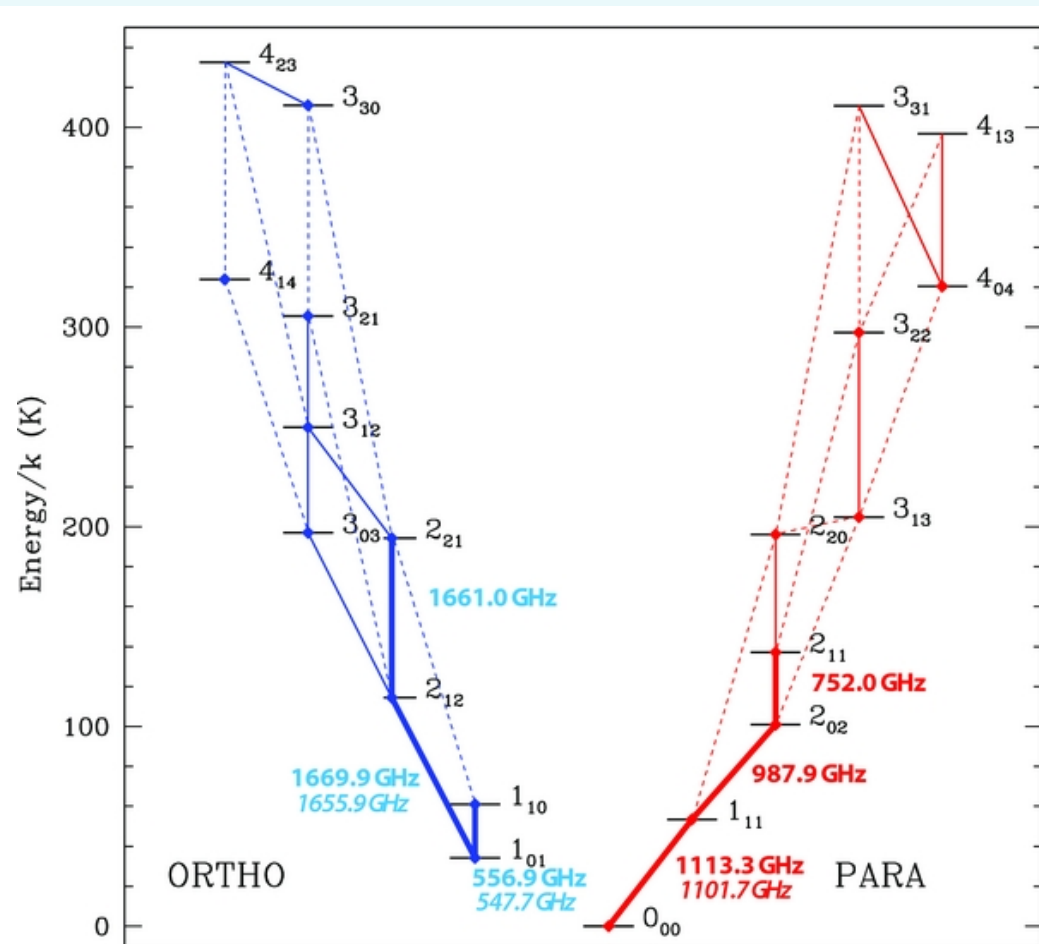
- Because H₂ has no dipole moment, J=1 → J=0 is not allowed.
- The lowest transition is: J=2 → J=0 (para-hydrogen)
- The next lowest is: J=3 → J=1 (ortho-hydrogen)
- $d = 7.414 \times 10^{-11}$ m (from Google); $\mu = 0.5 m_H$

$$E_J = J(J+1) \frac{\hbar^2}{2\mu d^2} = J(J+1) B$$

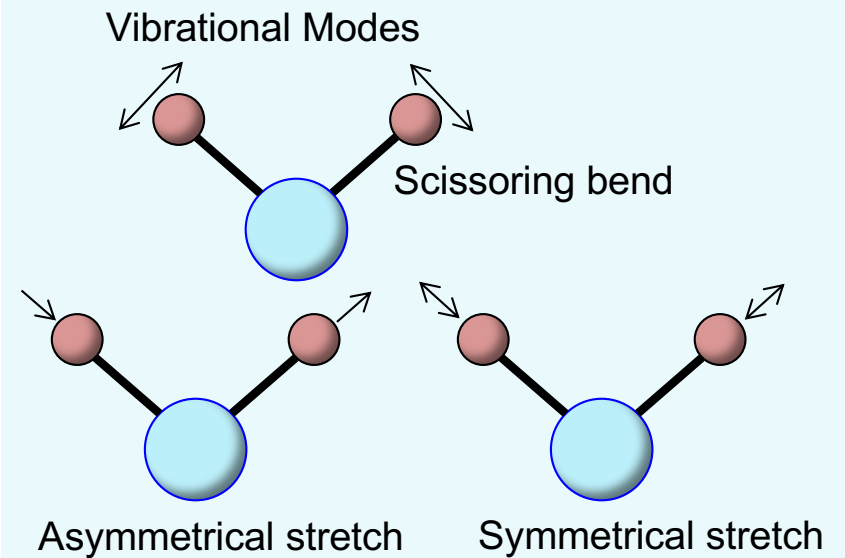
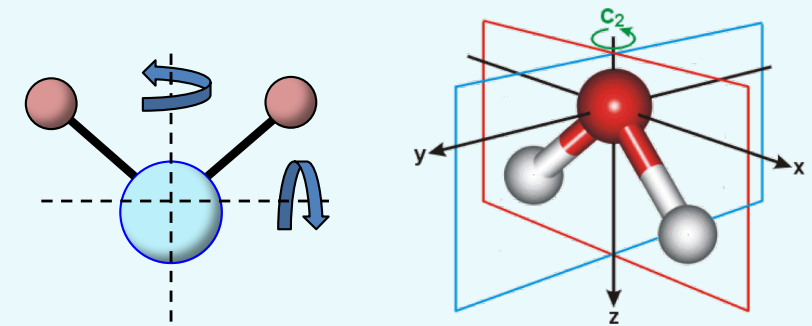
$$\Rightarrow \lambda(2-0) = 28 \mu\text{m}; \lambda(3-1) = 17 \mu\text{m}$$

- Similarly for hydrides, XH, where X=C, N, O, ... ; the ground state lines lie in the far infrared.
 - Estimate: $d \sim 10^{-10}$ m; $\mu = A_X/(A_X+1) m_H \sim m_H \Rightarrow B \sim 500$ GHz h
 - Lowest transition at 2B: $\lambda(1-0) \sim 300 \mu\text{m} [\mu/m_H] [d/\text{\AA}]^2$

- H₂O Cooling: H₂O quantum numbers J, K_{+1}, K_{-1} , $\Delta J = \pm 1$, $\Delta K = \pm 1, \pm 3$
 - If K_{+1}, K_{-1} have same parity => “para” state, otherwise “ortho” state



Credit: Flagey, et al. 2013, ApJ 762, 11

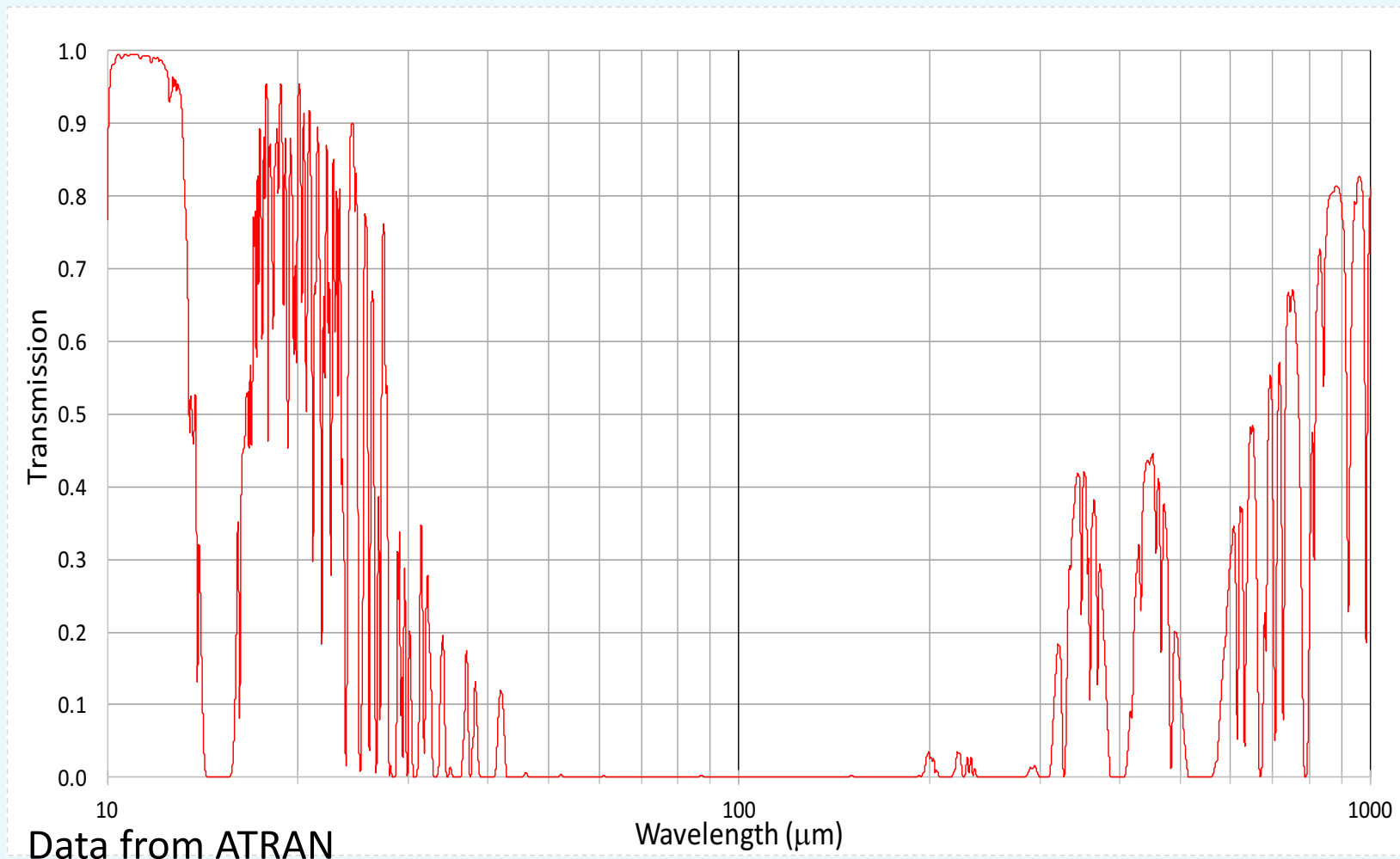




Water vapor makes the atmosphere opaque in the Far IR



- Good night at ALMA $\sim 700 \mu\text{m}$ water vapor column

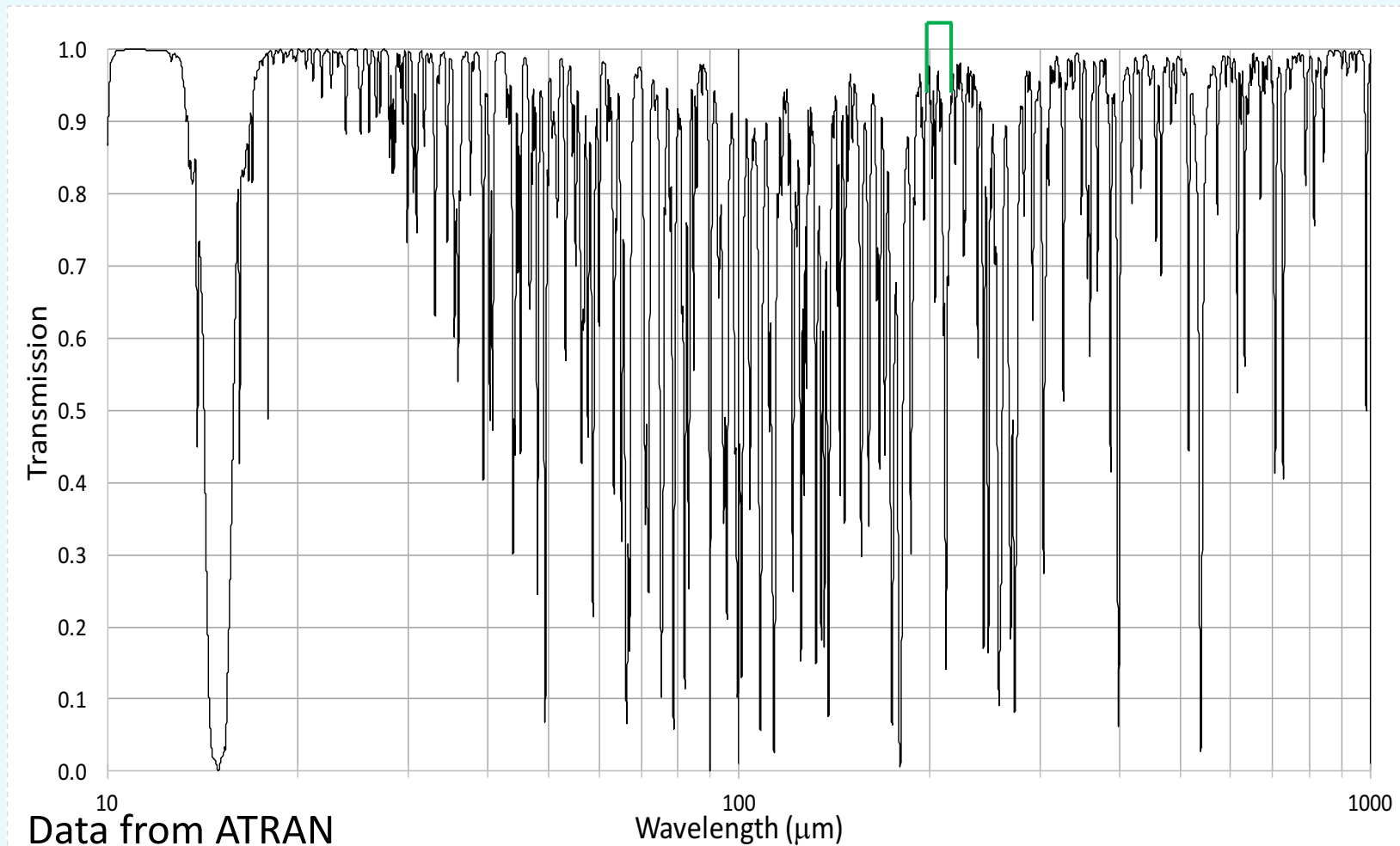




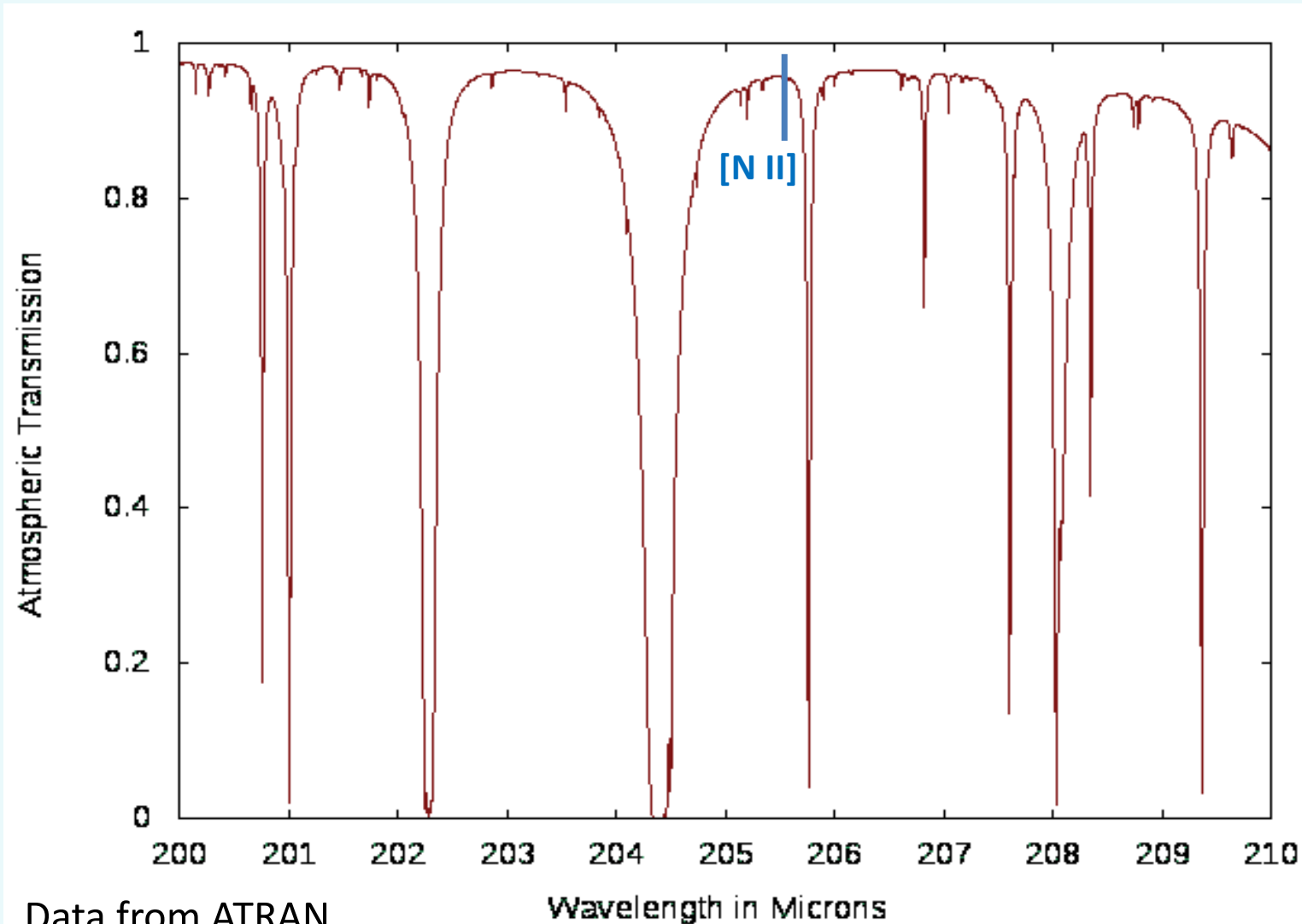
Water vapor in the atmosphere



- Good night on SOFIA ~10 μm water vapor

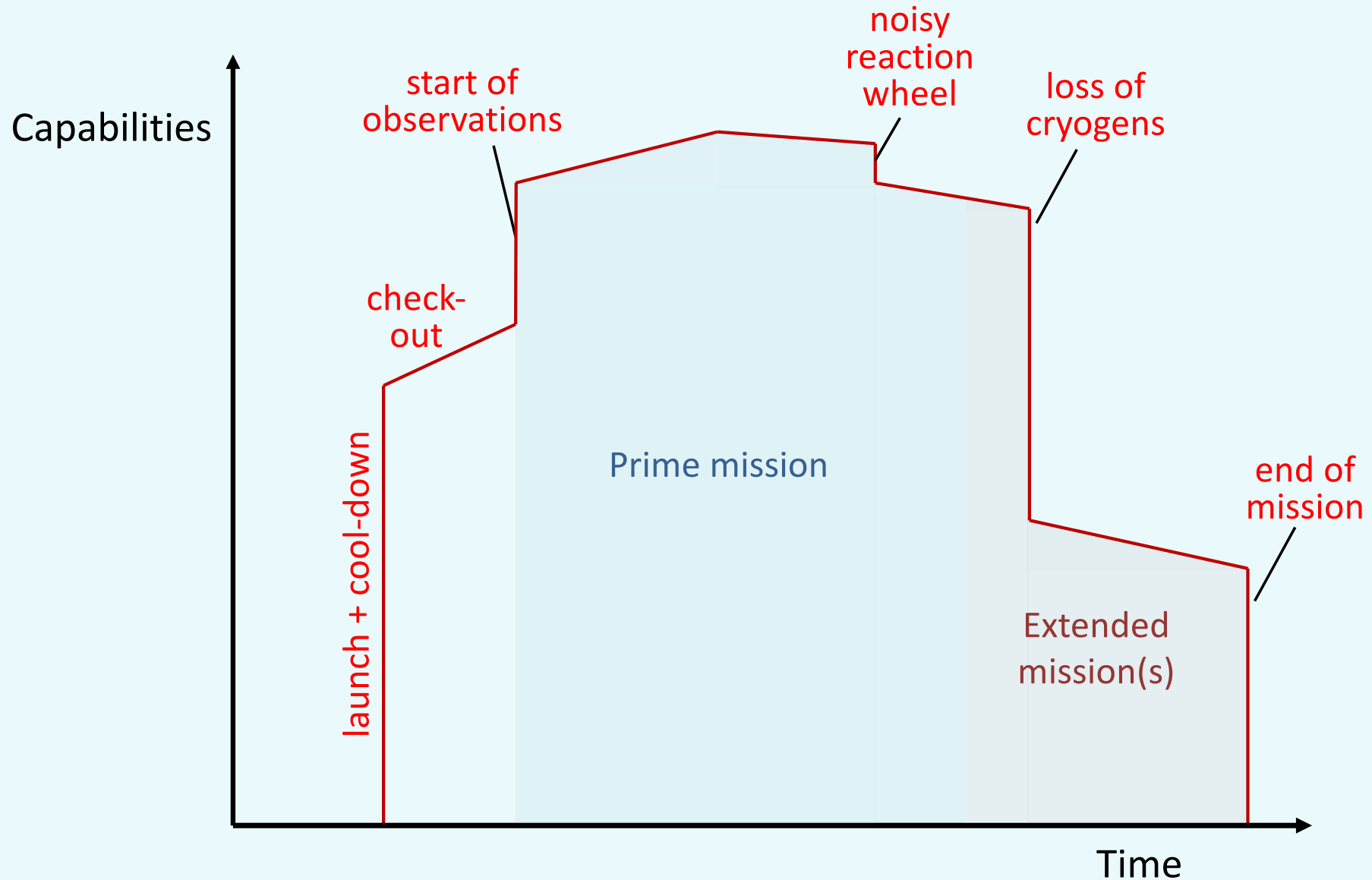


- Good night on SOFIA ~10 μm water vapor





The evolution of capabilities of space missions



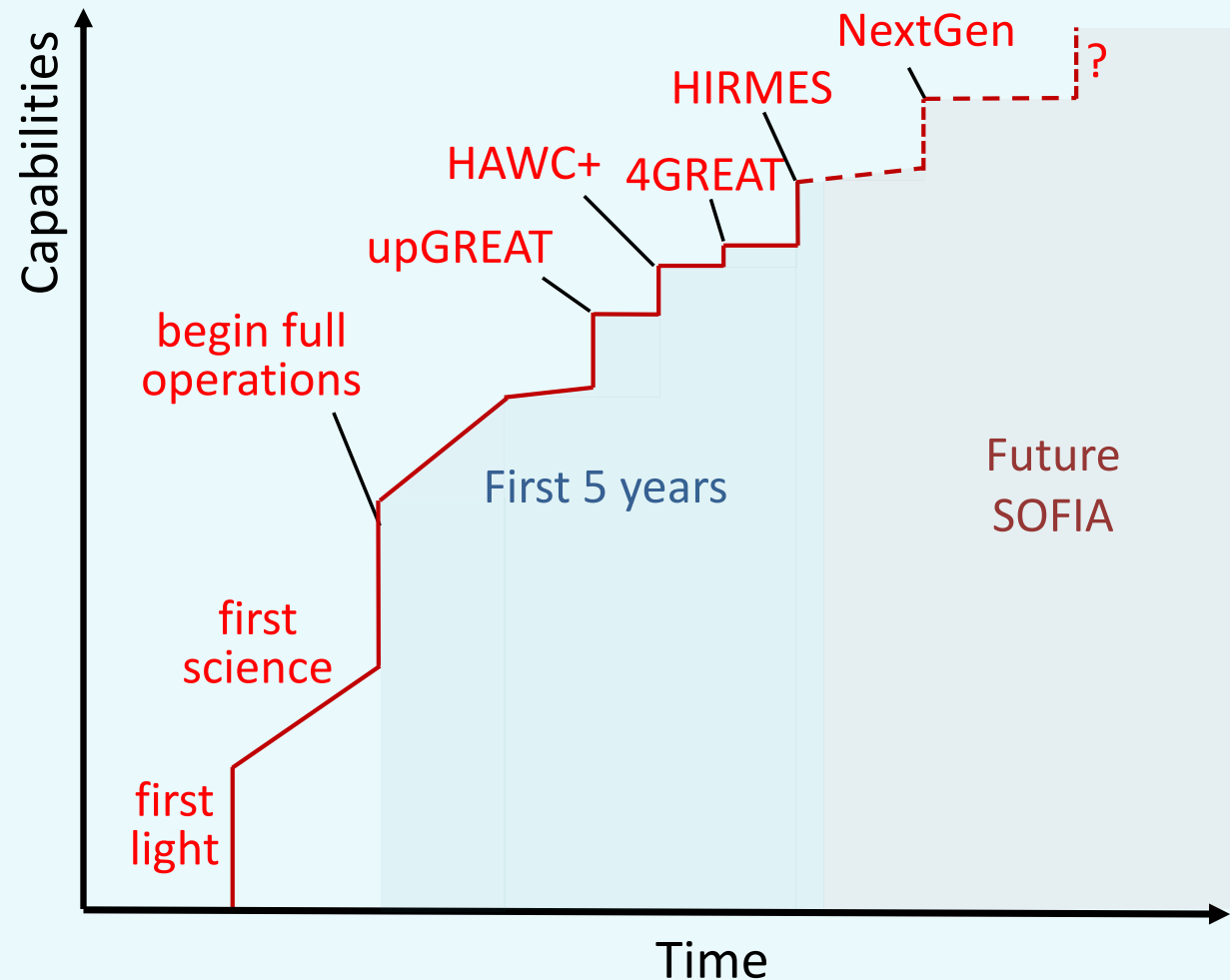


The evolution of **SOFIA** capabilities



SOFIA is not a space mission

- Hardware repairs & updates are possible on a relatively short time scale
- New instruments can be added to address current relevant science questions
- **SOFIA** tomorrow is much more capable than **SOFIA** today which is more capable than **SOFIA** yesterday





SOFIA is unique...



- SOFIA is the only observatory with access to infrared wavelengths above 99% Earth's water vapor
 - Its unique capabilities are necessary for quantitative answers to key science questions.
 - Through the development of new instruments, new science themes can be quickly addressed.
- SOFIA is an outstanding on-sky laboratory for development of future space-based instruments
 - Ample power, weight, 4K cooling, and computing facilities as well as personnel on board allow lower TRL instruments to be designed, built, flown, debugged, used, and perfected before adapting and qualifying for space.





SOFIA is unique...



- SOFIA maintains a working far-infrared community in preparation for a future space mission in the 2030s.
 - Need to maintain experts in dust physics, far- infrared spectroscopy and heating/cooling of gas and dust.
 - Far-infrared detectors are not commercially available; need technologists to continue advancing this field.



<http://www.sofia.usra.edu>

